### TEXAS ASM UNIVERSITY



No. 404 May 1992

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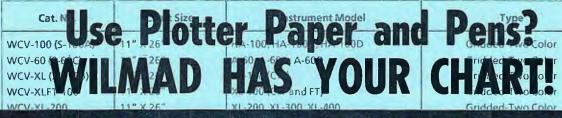
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### FORTHCOMING NMR MEETINGS

Gordon Research Conference: Magnetic Resonance in Biology and Medicine, Tilton, NH, July 13 - 17, 1992; Contact: Dr. A. M. Cruickshank, Gordon Research Center. University of Rhode Isalnd, Kingston, RI, 02881-0801; (401) 783-4011/3372; FAX: (401) 783-7644.

ISMAR 92 (The XIth Meeting of the International Society for Magnetic Resonance), Vancouver, B.C., Canada, July 19 - 24, 1992; Chairman: C. Fyfe. Contact: ISMAR 92, Dept. of Chemistry, University of British Columbia, Vancouver, BC, Canada V6T 1Z1. Tel: (604) 822-2293; FAX: (604) 822-2847; EMAIL: ismar@unixg.ubc.ca; BITNET: ismar@ubcmtsg.bitnet.

Science Innovation '92, New Techniques and Instruments in Biomedical Research (sponsored by the AAAS), San Francisco, July 21-25, 1992; Workshops on biomedical imaging. chemical and structural NMR; Plenary session on NMR on Fri., July 24. Contact: Science Innovation '92, P.O. Box 630285, Baltimore, MD 21263; fax (202) 289-4021.

34th Rocky Mountain Conference on Analytical Chemistry, Denver, Colorado, August 8-14, 1992; Contact: M. C. Goldberg, P.O. Box 25046 MS 424, Lakewood, CO 80225; (303)236-4728.

Eleventh Annual Scientific Meeting and Exhibition, Society of Magnetic Resonance in Medicine, Berlin, Germany, August 8-14, 1992; Contact: S.M.R.M., 1918 University Ave., Suite 3C, Berkeley, CA 94704; (415) 841-1899, FAX: (415) 841-2340.

XV International Conference on Magnetic Resonance in Biological Systems, Jerusalem, Israel, August 16 - 21, 1992; Contact: Prof. Gil Navon, XV ICMRBS, P. O. Box 50006, Tel Aviv 61500, Israel.; Tel. (972-3) 5174571, Fax: (972-3) 655674/660325.

MRI in the Applied Sciences, Duke University, Durham, North Carolina, October 25-28, 1992; Contact: Society of Magnetic Resonmance in Medicine, 1918 University Ave., Suite 3C, Berkeley, CA 94704; (510) 841-1899; FAX: (510) 841-2340.

High Resolution NMR Spectroscopy (a residential school), University of Sheffield, England, April 1993[sic]; Organizer: Dr. B. E. Mann (Sheffield); For information, contact Ms. L. Hart, The Royal Society of Chemistry, Burlington House, Piccadilly, London W1V 0BN, England; Tel.: 071-437-8656. Additional listings of meetings, etc., are invited.

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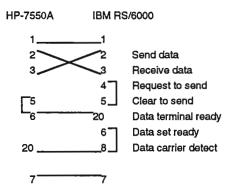
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Professor Barry L. Shapiro TAMU NMR Newsletter 966 Elsinore Court Palo Alto, CA 94303 April 10, 1992 (received 4/13/92)

### RE: Connecting an AM spectrometer and an IBM RS/6000 computer to a HP 7550 plotter.

#### Dear Barry.

We plot our Bruker AM-500 data on an HP-7550A plotter, and we can also plot data from an IBM RS/6000 computer that we use for off-line data processing. At the moment we are experimenting with Hare's Felix program Version 2.0. In order to eliminate any need to change either the configuration of the plotter or to change the RS-232 cables, the RS-232 cables from both the spectrometer and the IBM computer are attached to a serial auto switch. The switch we are using requires hardwire handshaking so the RS/6000 had to be configured for the DTR protocol. This auto switch can be set to 1 of 4 ports manually or it can rapidly scan its input ports and if the DTR line (20) of a port becomes active it locks on that port. In order for the automatic switch to work with the Risc system 6000 an RS-232 cable had to be wired as shown below:



The pref31 file has to be changed to ESC. P3: ESC. Y: IN; for hardwire handshaking. Felix adds a command to the start of the plot file that is PS3 or 4, where PS stands for page size. I am not certain if this is an old HPGL command or one that has been revived. Regardless it is not recognized by the HP 7550 plotter. Fortunately, the plotter ignores the instruction leaving only a complaint on the lcd display.

To facilitate transfer of data from the ASPECT computer for remote processing, I have written a program that transfers ASCII and data files between a PC and the ASPECT 3000. This program uses the same hardware as described by Markley et al. The data can then be transferred via ethernet to the RS/6000. There it is converted into the old format for Felix. If anyone has a listing of the locations of the parameters that can go into the new format, I would appreciate seeing it. The manual for 2.0 only lists a few parameters and there is an error in my copy of the manual. I can be reached by e-mail at ccottrel@magnus.acs.ohio-state.edu or the shorter address of cottrell.2@osu.edu. Please credit this to Alan Marshall's account.

Sincerely,

Charles E. Cottrell Senior Research Associate

Charles Cottul

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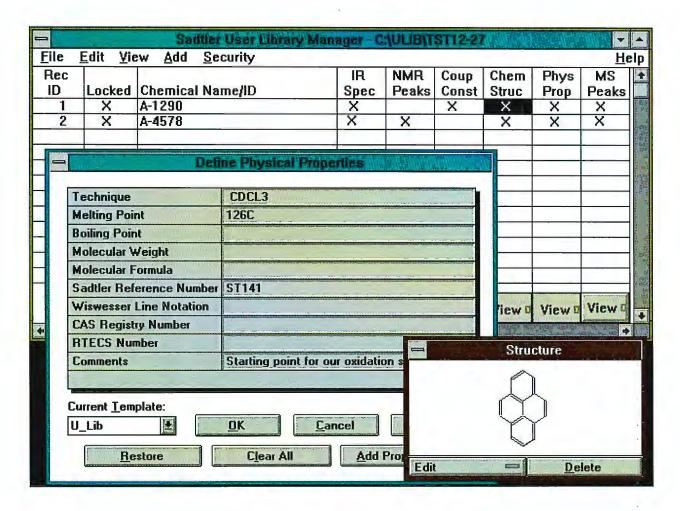
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March 28th 1992 (received 4/11/92)

Professor B.L. Shapiro, TAMU NMR Newsletter, 966 Elsinore Court, Palo Alto, CA 94303

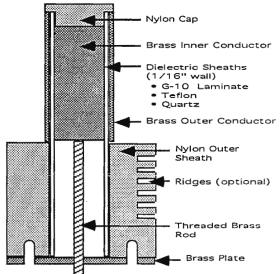
### High Voltage Variable Capacitor Design

Dear Professor Shapiro.

Thank you for your ultimatum. We have recently been designing and building radiofrequency coils for proton and phosphorus spectroscopy of the human brain in a 1.5 Tesla magnet. Some of the designs require high voltage variable capacitors in the range of 1-50pF which can withstand transmitter powers of up to 2 kW. Since commercially available capacitors in the desired capacitance/breakdown voltage/Q factor range are often (a) prohibitively expensive, (b) unavailable, (c) non-MR compatible (i.e. paramagnetic), (d) geometrically incompatible with the probe design, and (e) any combination of (a)-(d), we have been investigating the possibility of constructing our own capacitors.

The maximum voltage across the capacitor depends on the power of the RF transmitter P and the Q of the system of which the capacitor is a part. Given that  $V_{Rp-p} = V_{avg} \sqrt{2} = \sqrt{2*P*R}$ , and that the maximum reactive voltage across the capacitor  $V_{max} = V_{50p-p}*Q$ , then an amplifier output of 2 kW, with  $R = 50 \Omega$  and Q = 100, will generate an expected  $V_{max}$  of around 45kV across the capacitor.

Tubing constructed of fiberglass epoxy laminate (G-10), Teflon (PTFE), and fused silica (quartz) were chosen as the dielectric materials because of their high dielectric strengths and low dissipation factors. Dielectric strengths and loss tangents were 28.5 kV/mm and 2x10<sup>-2</sup>, 25 kV/mm and 2x10<sup>-4</sup>, and 35kV/mm and 5x10<sup>-4</sup> respectively. A cylindrical geometry was selected based on available materials and ease of construction. For a cylindrical capacitor with dielectric at a radius of ro, dielectric thickness d, length L, and relative dielectric constant  $\varepsilon_{\rm I}$ , capacitance is C =  $2\pi r_0 L \varepsilon_{\rm I} \varepsilon_0/d$ . Of the materials tested, quartz gave the highest Q factor but was prone to flashover due to lack of a tight seal between the quartz and the insulator separating the inner and outer conductors. The G-10 was very lossy giving unacceptably low O values. Teflon was the easiest to machine to avoid flashover, and when loaded with a sample, coil Q values were not significantly different between Teflon and quartz. Higher breakdown voltages should be obtainable by increasing the dielectric wall thickness.



The capacitor design illustrated here is relatively simple and can be built at low cost. Teflon appears to be the most suitable dielectric material for this application.

Sincerely yours,

Brian J. Soher

Irving Weinberg

Peter B. Barker

Peter Barner.



## Université de Lausanne - Faculté des Sciences INSTITUT DE CHIMIE MINÉRALE ET ANALYTIQUE

Place du Château 3, CH - 1005 LAUSANNE (Switzerland) tel. (021) 316 38 00 fax. (021) 316 38 01

Prof. Bernard L.Shapiro 966 Elsinore Court Palo Alto, CA 94303 U S A

15 April, 1992 (received 4/20/92)

### A RELATIVELY CHEAP AND RELIABLE METHOD TO STORE NMR DATA

Dear Prof. Shapiro,

For a long time storage of NMR data was a big problem in our Institute. Studying ligand exchange of metal complexes we use quite often the so called "isotopic substitution technique" to measure very slow rates. We mix for example a deuterated complex with a protonated ligand in a diluent and record <sup>1</sup>H spectra in fixed time intervals. Between 30 and 100 FIDs are then recorded and stored on disk. If we study temperature dependence and the dependence on ligand concentration, several hundred spectra have to be stored and kept until complete analysis. We are using Bruker spectrometers that are equipped with ASPECT computers. For a long time we transferred data for backup and storage to an ASPECT satellite station with a magnetic tape unit connected. This procedure was very time consuming, especially if you would like to read back an old spectrum sitting at the end of the tape reel.

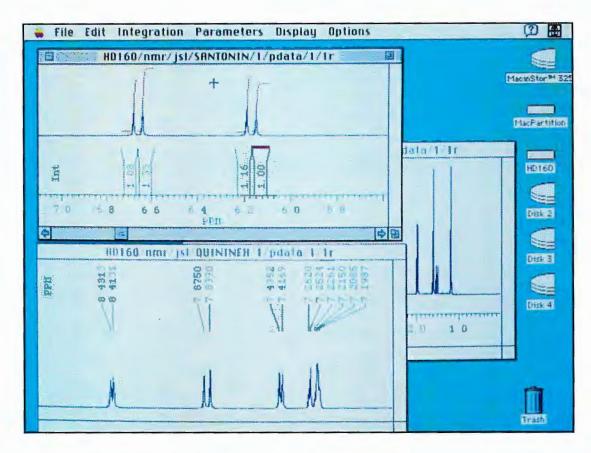
Since several months we own a REO-650 5.25" erasable optical disk system connected to a SUN workstation. At the beginning of this year we started to transfer data to the workstation via Bruknet (Brukers Ethernet protocol) and store it on optical disks. Care has to be taken to use the "P" option on the Bruknet demon running on the workstation. This option splits every datafile in two files: one for the parameters and one for the actual data. If the NMR data are transferred back to the ASPECT computer, Bruknet assembles the files in one, which can be read by the DISNMR program. The advantages of this method are:

- 1. fast storage and recovering of data because the access to individual files is easy
- relatively cheap and reliable storage medium (one optical disk can hold about 2 times 300
  Mbyte and costs about \$350, making this medium less expensive per Mbyte than magnetic
  tapes)
- because our PCs are linked to the workstation data can be transferred via Ethernet (ftp) to the PCs and treated there using Brukers WIN NMR

Dr. L. Helm

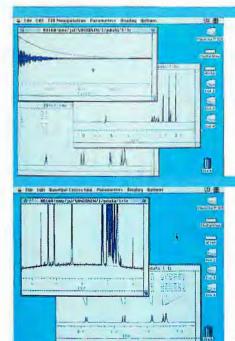
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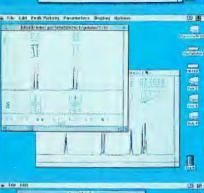


WIN-NMR can process FID and spectral data from all Bruker NMR spectrometers. The WIN-NMR data format used is identical to that of the UXNMR and UXNMR/P software. Data conversion routines are provided for all Aspect 2000/3000 based NMR spectrometers.

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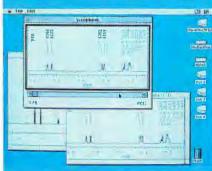
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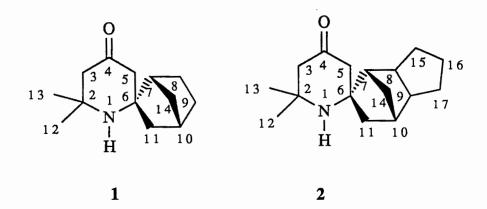
April 18, 1992 (received 4/22/92)

Dr. Bernard L. Shapiro Tamu NMR Newsletter 966 Elsinore Court Palo Alto, CA 94303

### Structure Analysis With 2D-NMR Technique

Dear Dr. Shapiro:

we have been interested in the synthesis of chiral nitroxides of the TEMPO type. The structures of two chiral precursors of some nitroxides are shown below. The structural problem is the orientation of the bridgehead CH<sub>2</sub> with respect to the piperidine ring, that is wether it is "in front" with respect to the nitrogen or "in back".



We have concluded that in both structures, The  $CH_2$  is back as drawn. The evidence for these structures is based upon INEPT,  $^1H_{-13}C$  COSY,  $^1H_{-1}H$  COSY and  $^1H_{-1}H$  NOESY spectra, as measured on a Bruker AC-270 NMR Spectrometer. The proton and carbon shifts are given in the attached Table.

The actual structures are based on NOESY enhancement observed between the bridgehead CH<sub>2</sub>-protons and those on carbons 5, 7, 10, and 11 for both structures. In addition, there is an enhancement between the protons on carbon 11 and those on the axial methyl group on carbon 2. In structure 1, it is clear that there is no NOESY enhancement between the protons on carbon 8 and 9 and those on carbon 5. This is less clear in structure 2 because one of the protons on carbon 8 or 9 is unpredictably far down field.





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Proton C	Chemica	l Shift	S	Carbon	n Chem	ical	Shifts		
Proton	1	2		Carbon	1	2			
H3 <sub>1</sub>		2.30	2.24		2			55.4	55.5
H3 <sub>2</sub>		2.27	2.28		3			54.1	55.0
H5 <sub>1</sub>		2.42	2.36		4			210.7	210.9
H5 <sub>2</sub>		2.29	2.41		5			49.2	54.0
H7		2.15	1.87		6			64.5	64.0
$H8_1$		1.42	2.48		7			47.0	51.2
$H8_2$		1.93			8			23.4	41.0
H9 <sub>1</sub>		1.23	1.88		9			28.8	40.7
$H9_2$		1.56			10			36.4	47.2
H10		2.20	1.84		11			49.2	48.6
H11 <sub>1</sub>		1.14	1.04		12			29.4	29.2
H11 <sub>2</sub>		1.70	1.69		13			32.5	32.6
H12a		1.19	1.18		14			38.0	32.4
H13e		1.27	1.24		15				31.8
H14 <sub>1</sub>		1.28	0.92		16				27.3
H14 <sub>2</sub>		1.58	1.86		17				31.4
H15 <sub>1</sub>			1.27						
H15 <sub>2</sub>			1.37						
H16 <sub>1</sub>			1.21						
H162	,		1.65						
H17 <sub>1</sub>	,		0.97						
H172			1.82						

Please credit this contribution to Dr. Thomas K. Leipert's account.

Sincerely,

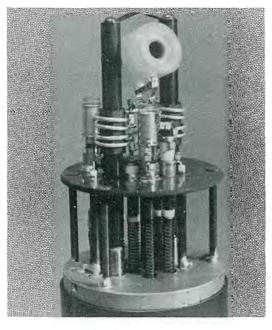
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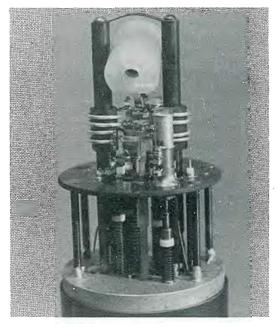


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> 1984	5 mm MAS at 9 kHz	> 1989	Precision Temperature Controller
> 1984	Multinuclear Observe for Solids	> 1990	Dynamic Angle Spinning
> 1985	19 mm MAS	> 1991	14 mm MAS at 7 kHz
> 1987	CRAMPS Probes	> 1992	3.5 mm MAS at 25 kHz
> 1987	Actively Shielded Gradient Probes	> 1992	300°C Pulsed Field Gradient Probe
> 1987	5 mm MAS at 17 kHz	> 1992	10 mm MAS at 10 kHz





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Dr. B. L. Shapiro, Editor TAMU NMR Newsletter 966 Elsinore Court Palo Alto, CA 94393 April 10, 1992 (received 4/16/92)

### A Field-Sweep CW NMR Using a Balanced Bridge/Network Analyzer

Dear Dr. Shapiro:

We use a Hewlett-Parkard 4195A Network/Spectrum Analyzer for probe tuning, and recently tested the device as a detector for CW NMR. Briefly, in the network analyzer mode, the HP4195A has an RF output (10Hz to 500MHz) and two inputs (one for reference port, one for test port); the amplitude and phase ratio of these two inputs can be measured as a function of frequency. With the addition of two reflection bridges (magic tee), a probe, and a swept magnetic field, an NMR is created. The magnet is normally associated with an MSL 200. The Kepco power supply is used to sweep the field about 22.4 kHz via the Z<sub>0</sub> shim.

Bridge circuits are common in old style NMR <sup>1,2</sup>. A coaxial bridge was used successfully as a nuclear quadruple resonance detector <sup>3</sup>, and lastly a balanced bridge system can also make low temperature NMR more convenient <sup>4</sup>.

The test sample is adamantane in a 7 mm rotor spinning at the magic angle at a rate of 3 kHz in a field of 4.7 Tesla. Figure 1 shows the actual configuration of the system. Briefly, the first bridge is used as a power splitter, the second is the balanced bridge. Under normal pulsed <sup>1</sup>H NMR condition, a five line pattern is expected with a really great signal-to-noise ratio. In the first attempt (shown in Figure 2), the S/N seems to be about 100-fold less with this setup. Improvements could come with better impedance balancing of the bridge and incorporation of digital lock-in amplification. A modified version of this configuration will serve as one of our future NMR detectors for the new-low temperature, high-magnetic field system.

Commands and suggestions will be appreciated.

Sincerely,

Vino W

Les Butler

<sup>1.</sup> P. Kesselring and M. Gautschi, J. of sci. Instrum., 1967, vol 44, p911-916

<sup>2.</sup> D. R. Torgeson, Rev. of Sci. Instrum., 1967, vol 38, p612-615

<sup>3.</sup> M. P. Klein and D. E. Phelps, Rev. of Sci. Instrum., 1967, vol 38, p1545-1546

<sup>4.</sup> T. J. Gramila, "Experimental Techniques and Special Devices", *Experimental Techniques in Condensed Matter Physics at Low Temperatures*, Addison - Wesley Publishing Company, Inc.

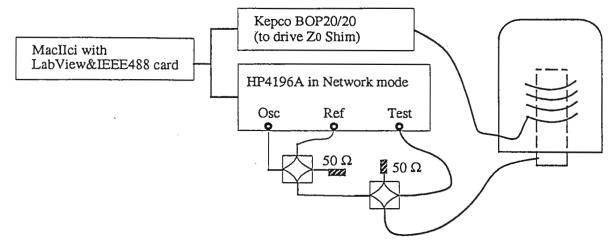


Figure 1. Experimental setup

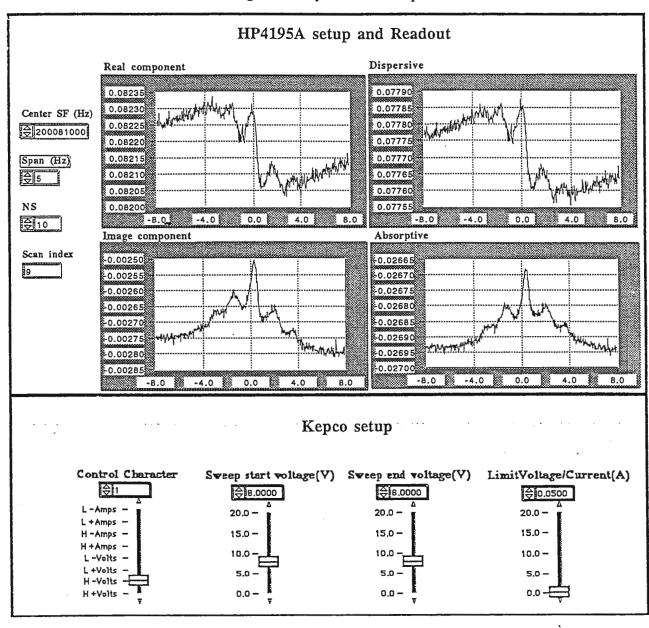


Figure 2. LabView program front panel

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### Chemagnetics Workshop Set For August 7 & 8

Dear NMR Colleagues:

I'd like to invite you to an NMR workshop that you won't want to miss -- especially if you're attending the Rocky Mountain Conference. Each year, following the conference in Denver, Chemagnetics invites a number of NMR scientists to give seminars on advanced topics and organizes a 2-day opportunity for spectroscopists to visit and get up-to-date on the latest in solid-state NMR. This year's workshop will be held August 7 and 8, 1992 in Fort Collins, Colorado.

We are privileged to have luminaries from both academia and industry. Below is a list of seminars that will be presented on Friday. Saturday will feature hands-on demonstrations of advanced solids techniques, including DOR, fast MAS, and liquids microimaging.

Speakers and topics for the Friday session:

- Dr. Brad Chmelka, UC Santa Barbara
- Dr. Jim Haw, Texas A&M University
- Dr. Dave Duff, Raychem Corporation
- Dr. Robert Botto, Argonne National Laboratories
- · Dr. Lucio Frydman, UC Berkeley
- · Dr. Robert Tycko, AT&T Bell Labs
- Dr. Ann McDermott, Columbia University
- Dr. Hellmut Eckert, UC Santa Barbara
- · Dr. Gerald S. Harbison, SUNY

Practical Applications of 2-D exchange NMR in Polymers

Recent Developments in MAS NMR

Triple Resonance Experiments on Protonated and Fluorinated Polymers

NMR Imaging: A Chemical Microscope for the Analysis of Heterogeneous Materials

NMR Imaging of Flow

New Approaches to Structure Determination in the Solid State

Applications of Rotational Resonance in Enzyme-active Sites

Heteronuclear X-Y Double Resonance NMR Experiments in Solid Inorganic Materials

Solid-State NMR of Highly-Oriented DNA Fibers and Drugs Bound to DNA

If you would like to attend, I urge you to register today. Registrations are limited to 60 people on the first day and 15 for day two. There is no fee to attend the seminar, but registrants are responsible for their own transportation from the Denver airport to Fort Collins (approx. \$13), for the hotel, and meals. To help defray your costs, rooms at Chemagnetics' discount rate of \$34 have been booked at the nearby Ramada Inn. Please call to register today at 1-800-4-OTSUKA and ask for Jan Anderson.

I look forward to seeing you.

Robert Wind

Director of Research and Development

JRobert A. Wind

### Postdoctoral Position

The Research and Development Department at Chemagnetics has a postdoctoral opening for a Ph.D. NMR spectroscopist. Experience with solids and NMR and/or MRI techniques and hardware is required. The work involves the development and evaluation of novel solid-state techniques in biosystems. The successful candidate must be interested in working in a problem-solving environment and be able to work as a member of a team.

Interested candidates should send their resume to: D. Hasler, Chemagnetics, Inc., 2555 Midpoint Drive, Fort Collins, CO 80525. For further information, contact Robert Wind (303) 484-0428. EOE



1222

Faculteit der Scheikundige Technologie en Materiaalkunde Technische Universiteit Delft

Laboratory of Organic Chemistry and Catalysis Delft University of Technology Julianalaan 136, 2628 BL Delft, The Netherlands

### Discrimination of Diastereomeric

### Dioxolane Compounds Using Gd(III)-Induced Carbon-13 Relaxation Rate Enhancements

(received 4/6/92)

Dear Dr. Shapiro,

A topic of our investigations is the addition of hydroxyl compounds to unsaturated carboxylic acids with the use of lanthanide(III) ions as template. A specific example is the addition of ethylene glycol to acetylene dicarboxylate (1), which yields a 2,2-disubstituted dioxolane. It appeared that this reaction proceeds via a diadduct, in which two ethylene glycol units are attached to the acetylene dicarboxylate chain, and which gives an intramolecular ring closure only after decomplexation. Addition of glycerol (2) to acetylene dicarboxylate yields a mixture of the two dioxolanes 3 and 4 (see scheme), of which one, further called A, is three times more abundant than the other (B). With the use of standard <sup>13</sup>C NMR, NOESY, and NOE experiments we could not clarify the configuration of these diastereomers.

Gd(III)-induced <sup>13</sup>C relaxation rate enhancement measurements are normally used to get an idea of the steric conformation of organic ligands in lanthanide complexes. We successfully applied this technique to discriminate between the two dioxolane diastereomers. From the data given in the table it can be concluded that in both diastereomers the Gd(III) ion is preferably coordinated by C6 and probably one of the ether oxygens as shown by a two times higher relaxation rate enhancement for C6 than for C8. The relaxation rate enhancement of C9 of A was only a third of that of C9 of B. This implicates that in the latter compound C9 is at the same side of the dioxolane ring as C6 and therefore structure 3 can be assigned to it.

By this example we have shown that Gd(III)-induced <sup>13</sup>C relaxation rate enhancement measurements can be a powerful tool in discriminating between strongly related compounds. Yours sincerely,

Jurriaan Huskens,

Joop A. Peters,

carbon	1/T <sub>1</sub> [:	10 <sup>4</sup> s <sup>-1</sup> ]
atom	A (4)	B (3)
C2	1.565	1.588
C4	0.627	0.910
C5	0.515	0.804
C6	5.286	5.310
C7	0.885	0.940
C8	2.499	2.450
C9	0.223	0.646

Gd(III)-induced  $^{13}C$  relaxation rate enhancements (100 MHz, 30 $^{\circ}C$ ) of 3 and 4 (2 M in  $D_2O$ ) extrapolated to  $\rho = [Gd]/[ligand] = 1$ 

Herman van Bekkum

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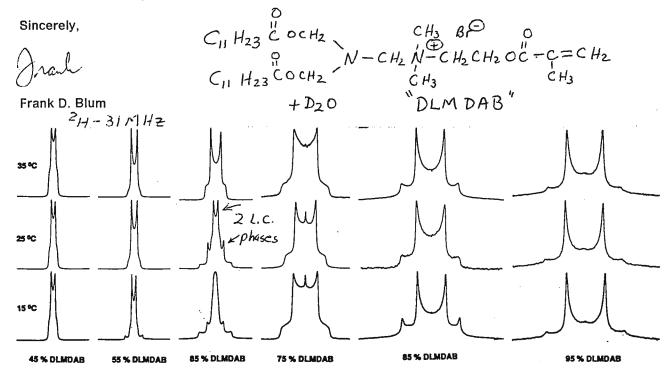
March 28, 1992 (received 3/31/92)

Dr. Barry Shapiro TAMU Newsletter 966 Elsinore Ct. Palo Alto, CA 94303

### Dear Barry:

As you can see, I am at IBM. I will be on leave here for a half year. I would like to share some work that Jack Uang did with me at Missouri. We have synthesized a new surfactant that has a very long name that I will call DLMDAB. This surfactant forms liquid crystals with water over a wide range of temperatures. The use of deuterium oxide with DLMDAB provides a convenient way of probing the phase behavior of such a system. We find that along with X-ray, a good deal of information can be obtained about the system. Shown below is the partial NMR phase diagram of the system.

At 15 and 25 degrees, two different liquid crystals are formed at high and low surfactant concentrations. At 15 degrees, there is an apparently isotropic phase between them. At 25 degrees there is an isotropic phase plus the presence of two superimposed liquid crystalline powder patterns. This is from two liquid crystalline phases in equilibrium. The only difference between these is the different amounts of water in each. This is the third system that I know of where this type of behavior has been seen. If anyone wants a preprint of the paper, let me know. (Please credit this to my account in Missouri.)

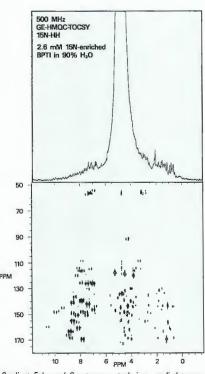


# Gradient Enhanced Spectroscopy: a new, practical answer

By Frank Huang, PhD and Paul Calderon, MS

Making Gradient Enhanced Spectroscopy (GES) a viable method for high resolution spectroscopy has long been of interest to researchers. The obvious benefits in speed and information content were too often overshadowed by the drawbacks of signal loss and distortion.

Technology developed at GE NMR Instruments has overcome these



Gradient Enhanced Spectroscopy technique applied to an HMOC-TOCSY experiment demonstrates excellent water suppression without the need for presaturation or selective excitation.

### Research Implications

- ► Speed without phase distortion or signal loss
- Practical for ID, 2D, 3I and 4D experiments.
- Accommodates proton and heteronuclear GES techniques.

challenges. The S-17 Gradient Enhanced Spectroscopy Accessory with integrated inverse probehead makes GES practical for a broad range of applications in 1D, 2D, 3D and 4D experiments.

### A better design

The use of a three-axis, actively-shielded gradient set allows GE to overcome the inherent draw-backs of previous GES technology—most notably, phase distortion and signal loss:

Active shielding. Eddy current effects are the major source of phase distortion in GES spectra. Active shielding prevents interaction of strong gradients with magnet and shim components—the source of eddy current effects.

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### Applications advantages

With its integral inverse probehead, the S-17 Accessory can accommodate proton as well as heteronuclear GES techniques.

Gradient fields of 35 G/cm are able to suppress water in aqueous samples and to improve performance in heteronuclear experiments. Other applications advantages:

- ► Eliminates phase cycle requirements and subtraction error.
- ► Reduces T1 noise.
- Reduces collection times for 2D, 3D and 4D data sets.
- Provides lineshape independent water suppression in multiple quantum coherence selection experiments.
- Provides lineshape independent water suppression via diffusion differences for large molecular weight samples.
- ► Improves water suppression in experiments using selective time reversal RF pulses.
- Separates cross-correlation and exchange phenomena in NOESY experiments.
- ➤ Distinguishes chemicals inside the cell from those outside in whole cell applications.

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Or call toll free:

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**GE NMR Instruments** 

### Gradient-Enhanced 15N HMQC

The pulse sequence and the coherence pathway diagram for a 15N GE-HMQC are shown in Fig. 1. The pulse sequence was a standard <sup>13</sup>C GE-HMQC experiment (l) with different gradient amplitudes to account for the difference between the gyromagnetic ratios of <sup>13</sup>C and <sup>15</sup>N. The 90° proton pulse creates transverse magnetization which evolves into an anti-phase state with respect to J(NH) coupling at the end of the period  $\triangle$  (where  $\triangle$  =  $1/(2I_{NH})$ ). The antiphase components are converted into heteronuclear zero- and double-quantum coherence by the <sup>15</sup>N 90° pulse and the multiple quantum coherences are allowed to evolve during t<sub>1</sub>. The 180° <sup>1</sup>H pulse in the center of the evolution period serves to eliminate the <sup>1</sup>H chemical shift evolution, yielding pure <sup>15</sup>N chemical shifts along that axis. The zero- and double-quantum signals are then coherence-order labeled by the gradient pulses G1 and G2. After conversion into antiphase proton magnetization by the last <sup>15</sup>N 90° pulse, the desired components are refocused by the gradient G3 and detected. The application of a gradient pulse results in a phase factor being applied to the magnetization which is dependent upon gradient strength, duration, the distance from the gradient isocenter, the gyromagnetic ratios of the coupled nuclei, and the desired coherence order. The relative amplitudes of the labeling and refocusing gradient pulses will determine the selection of a specific coherence pathway and are calculated to suppress magnetization components arising from the solvent and other protons not coupled to 15N spins.

The fundamental principle of coherence selection using gradients is that for a pathway to be detected, the cumulative phase factor during the acquisition must be zero:

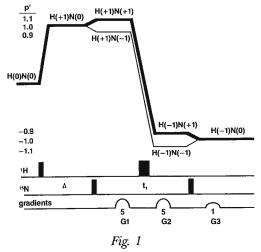
$$G_1 P_i + G_2 p_2' + G_3 p_3' = 0.$$
 [1]

The subscripts denote steps in the pulse sequence where p' defines a composite coherence order for the heteronuclear case which includes the gyromagnetic ratios of the coupled nuclei:

$$p' = p_1 H + (r_{15} N / r_1 H) p_{15} N$$
 [2]

and <sup>p1</sup>H and <sup>p15</sup>N are the coherence orders for the <sup>1</sup>H and <sup>15</sup>N spins respectively.

In the coherence pathway diagram, the relevant values of p' are given to the left and the relative gradient areas



Pulse Sequence and the coherence pathway diagram for a  $^{15}N$  GE-HMQC experiment.

(gradient strength x duration) are given next to each gradient pulse. The following pathway (shown in Fig. 1):

$$H(+1) \rightarrow H(+1)N(0) \rightarrow H(+1)N(+1) \rightarrow H(-1)N(+1)$$
  
 $\rightarrow H(-1)N(0)$ 

is detected using a 5:5:1 ratio of gradient areas, since according to Equation 2:

$$5(1.1) + 5(-0.9) + 1(-1.0) = 0$$
 [3]

where the numbers in the parentheses refer to the composite coherence orders. Using these relative gradient areas, protons not coupled with <sup>15</sup>N spins may pass through an alternate pathway:

$$H(+1) \rightarrow H(+1) \rightarrow H(+1) \rightarrow H(-1) \rightarrow H(-1)$$

which results in a net phase factor:

$$5(1.0) - 5(-1.0) + 1(-1.0) = -1$$
 [4]

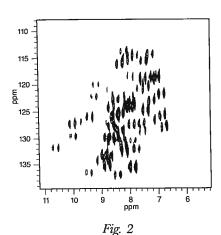
Thus, signals from this pathway remain defocused during the acquisition.

A 2D <sup>15</sup>N GE-HMQC spectrum of <sup>15</sup>N enriched BPTI is shown in *Fig.* 2. The spectrum was collected using a 5 mm inverse probe on an Omega<sup>TM</sup> PSG 500 spectrometer equipped with an S-17 gradient accessory. Half-sinusoid-shaped gradient pulses were applied simultaneously along the X, Y and Z axes with a maximum gradient strength of  $\approx$ 20 Gauss/cm and a duration of 3.5 ms. A matrix size of 2048 × 128 resulted in 3.5 Hz resolution in the  $\omega_2$  dimension and 10 Hz in the  $\omega_1$  dimension. No decoupling was applied.

Gradient-enhanced experiments provide a viable alternative to traditional phase-cycling methods for the selection of coherence pathways. In cases where the sensitivity is adequate, gradient selection can substantially reduce the collection time in multi-dimensional experiments. The <sup>15</sup>N GE-HMQC data presented here has none of the t<sub>1</sub>-noise from cancellation artifacts usually present in phase-cycled versions of the HMQC experiment. In addition, since the suppression of the single-quantum signals is done prior to acquisition, the receiver gain may be increased, which results in a substantial increase in signal-to-noise. For these reasons, gradient pulses should be the method of choice for coherence selection in HMQC experiments.

#### Reference

1. R.E. Hurd and B.K. John, J. Magn. Reson. 91, 648 (1991).



A 2D GE-HMQC spectrum of  $^{15}N$  enriched BPTI. The sample was 2.6mM in 90%  $H_2O$ . The data collection time was 2.6 hours.



April 2, 1992 (received 4/7/92)

Dept. of Inorganic Chemistry

# First Observation of Rhodium—Rhodium Two—Bond Spin—Spin Coupling.

Dear Prof. Shapiro,

During the last decade,  $^{103}$ Rh-NMR has grown considerably in importance as a tool for rhodium chemists despite the very low relative sensitivity ( $D_p = 3 \cdot 10^{-5}$ ) and long relaxation times.

Our group is presently interested in structural investigations in aqueous solution of the hydroxo bridged polymeric species formed on hydrolysis of rhodium(III). Due to their kinetic stability several species can be isolated one by one from a hydrolysed mixture by ion—exchange chromatography and studied at room temperature. <sup>103</sup>Rh—NMR has been the method of choice; however, we have also measured <sup>17</sup>O—NMR, <sup>1</sup>H—NMR, EXAFS, LAXS, *far*—IR and Raman spectra.

Beside the monomer,  $[Rh(OH_2)_6]^{3+}$ ,  $\delta_{Rh} = 9915.8$  ppm, and the dimer,  $[(H_2O)_4Rh(\mu-OH)_2Rh(OH_2)_4]^{4+}$ ,  $\delta_{Rh} = 9997.4$  ppm, a third species was obtained which gave two  $^{103}Rh-NMR$  signals having relative intensities 2:1,  $\delta_{Rh} = 9966.4$  and 10003.6 ppm, respectively. This result surprised us since in the otherwise analagous chromium(III) system the trimeric species has been found to have a triangular arrangement of equivalent chromium atoms. A high resolution spectrum was recorded of the suspected trimer and after Gaussian resolution enhancement the signals at 9966.4 and 10003.6 ppm were found to be split into a doublet (1:1) and a triplet (1:2:1) respectively,  $^2J_{Rh,Rh} = 1.5$  Hz. The observed spin—spin coupling was interpreted in terms of an AX<sub>2</sub> spin system ( $I_{Rh} = \frac{1}{2}$ ): the triplet arises from the central nucleus (A) through coupling to the two equivalent terminal rhodium nuclei (X), which in turn give rise to a doublet.

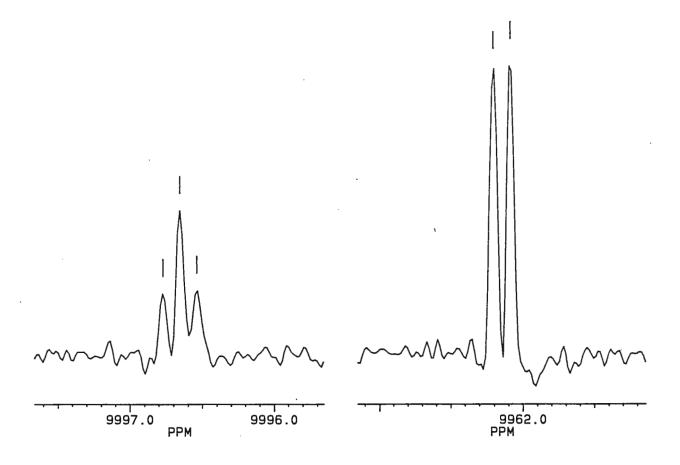
The splitting pattern and the ratio in the observed intensities (A:X = 1:2) between the two signals allowed us to confirm both the nuclearity and the structure of the trimer,  $[H_2O)_4Rh(\mu-OH)_2Rh(OH_2)_2(\mu-OH)_2Rh(OH_2)_4]^{5+}$ , from a single experiment. To our knowledge this is the first observation of rhodium-rhodium two-bond spin-spin coupling.

Yours sincerely,

Michael Read

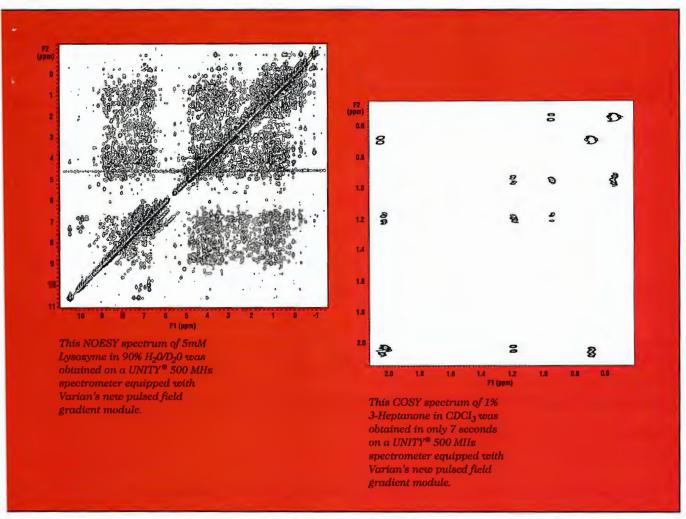
Julius Glaser

Magnus Sandström



Expansion of the resolution–enhanced  $^{103}\text{Rh}$ –NMR spectrum of an aqueous solution ([Rh] = 0.4 M) of the linear trimer, [H<sub>2</sub>O)<sub>4</sub>Rh( $\mu$ –OH)<sub>2</sub>Rh(OH<sub>2</sub>)<sub>2</sub>( $\mu$ –OH)<sub>2</sub>Rh(OH<sub>2</sub>)<sub>4</sub>]<sup>5+</sup>, recorded at 12.6 MHz, 298 K. The spectrum is referenced to  $\Xi(^{103}\text{Rh}) = 3.16$  MHz.

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DIVISION OF CHEMISTRY AND CHEMICAL ENGINEERING GATES AND CRELLIN LABORATORIES OF CHEMISTRY

Dr. Bernard L. Shapiro TAMU NMR Newsletter 966 Elsinore Court Palo Alto, CA 94303

JOHN D. ROBERTS INSTITUTE PROFESSOR OF CHEMISTRY, EMERITUS

April 13, 1992 (received 4/15/92)

Dear Barry: Proton Shifts in Three-Center, Two-Electron C-H-C Bonds in Carbocations

Tipping of your orange spot (it was black for RLS) causes me to call the attention of the TAMU fraternity to the unusual proton shifts reported for three-center, two-electron C-H-C bonds in carbocations. The situation is very nicely summarized by Cornell chemists, J.E. McMurry and T. Lectra in Accounts of Chemical Research, 1992, 25, 47-53. Because of prior entanglements in the carbocation field, I have been much intrigued by these shifts of the bridging protons, which can be observed for bicyclo[4.4.4]tetradecyl and

The bicyclo[4.4.4]tetradecyl carbocation

related cations in magic acid solutions, often even at room temperature. The facts are that this kind of proton comes at -3.5 ppm in the [4.4.4] cation, where the C-H-C angle is expected to be 180° and moves still farther upfield when the C-H-C angle is less than 180° and reaches -6.5 ppm with the [6.4.2] cation, where the angle is estimated to be 113°. An essentially linear correlation exists between the C-H-C angles and the proton shifts. McMurry states that the reasons for the correlation of proton chemical shift and C-H-C angle are not clear. And it does not seem that the shifts are consistent with simple considerations of the electron densities surrounding the bridging protons.

The large diamagnetic shifts could result from the second-order paramagnetic effect operating at the carbons so as to produce a large upfield shift of the bridging protons in something like the way that the nucleus of a non-hydrogen-bonded hydrogen in an O-H group has a large upfield shift. Thus, the primary shift effect arises not from the electron density at the protons, but from the circulation of electrons around the bridge carbons, when the C---C axis is perpendicular to the applied field. This circulation would produce a paramagnetic effect at the carbons and a consequent diamagnetic effect at the proton. In order to make sense of the changes of shift with the C-H-C angles, I have to postulate that this effect on the shift for the bridging proton in the [4.4.4] cation is larger than observed, perhaps -7 ppm, or even more. The effect might be enhanced because the carbon valence shells of the C-H-C groupings are not fully occupied, and so make the paramagnetic circulations larger than for normal saturated carbons.

Two effects might be expected to operate when the C-H-C angle becomes smaller than 180°. First, when the C---C axis is perpendicular to the magnetic field, the proton is not expected to be in the same favorable position to experience the maximum diamagnetic effect from the circulation around the carbons relative to when the C-H-C angle is 180°. This should cause paramagnetic shifts, contrary to what is observed. However, an opposite effect would result from contributions of a second-order paramagnetic effect directly at the bridging hydrogen as the result of admixture of some p-character to the hydrogen bonding orbitals by the electrostatic influence of the positive carbons flanking it. This sort of effect has been invoked to account for the substantial paramagnetic proton shifts observed for linear O-H--O hydrogen bonds. Any such contribution to the C-H-C proton shifts would surely be expected to diminish when the C-H-C bond angle bnecomes less than 180°. Thus, if this effect is important, but not important enough to completely overcome the diamagnetic shift, then one can rationalize the shifts. I expect that the intrepid readers of TAMU NMR will have superior explanations based on gauge invariance calculations, but how about something qualitative that we can use to help students understand shifts better?

With best wishes,

Very truly yours,

### LABORATOIRE DE RÉSONANCE MAGNÉTIQUE NUCLÉAIRE UNIVERSITÉ CLAUDE BERNARD LYON I



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March 30, 1992 (received 4/6/92)

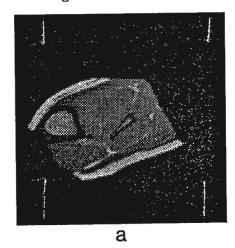
Dr Bernard L. SHAPIRO TAMU NMR Newsletter 966 Elsinore Court Palo Alto, CA 94303

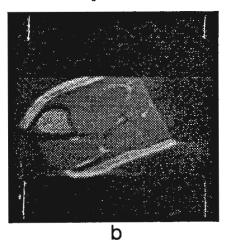
Cutting ham in the k space

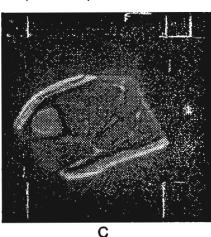
Dear Dr Shapiro,

It is well known that imaging can be performed using the half of the k space. This propriety works well in MRI as in Optics and it may be turned into profit in order to limit the amount of data and to reduce the time necessary for an image acquisition. This feature may considerably accelerate 3D imaging or any kind of imaging performed with low sensitivity signals of porous materials or polymers .

This fact is illustrated using a 2D experiment on ham at .13 Tesla where full echoes are acquired in the time domain as well as in the coding gradient direction. The images are reconstructed in three manners. First, the normal 2DFT is performed (Fig.a). In Fig.b, the k space is cut orthogonally to the time direction at the top of the echoes. A first order phase correction is performed after FT along the time domain. This phase correction may be required because sampling does not necessarily coincide with the top of the echo. It is also possible to cut the k space orthogonally to the coding gradient direction (Fig.c), in this case, the top of the echo is perfectly defined when the phase encoding is zero. For cases b and c, the signal was zero filled in order to keep the same matrix size (128x128).







5 mm slice of Italian Parma ham with bone

We observe that signal to noise ratio is slightly reduced. In case b, the image quality seems to depend on the phase correction. Nevertheless, in all cases the image quality is not fundamentally different, and the use of half k space is efficient.

Sincerely,

Y. CREMILLIEUX

hy

A. DIOP

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Dr. Bernard L Shapiro 966 Elsinore Court Palo Alto, CA 94303

#### RESOURCE-LIMITED OXYGEN SHIFTS?

March 11th, 1992 (received 3/30/92)

Dear Barry,

Anyone who measures <sup>17</sup>O NMR shifts of transition-metal oxyanions will know that some peaks have an annoying way of disappearing, usually behind more dominant resonances. It can be hard work to be sure of their presence, e.g. by pH variation or careful integration. One of the problems is that one is often unsure of the likely shift of the missing resonance. I would like to suggest a simple rule which we find helpful in such cases. After compiling a fair amount of data, we have noticed that the sum of the oxygen shifts divided by the number of metal atoms is roughly constant for a given metal, whatever the geometry of the oxyanion. The table below gives this shift sum per metal for known isopolyvanadates, molybdates and tungstates. It also gives a few sums for heteropoly species with central, tetrahedral XO<sub>4</sub> units. These are calculated by ignoring the shifts of the central oxygens; in fact they never differ much from those of the isolated XO<sub>4</sub><sup>n</sup> unit. The shift sums do vary somewhat from metal to metal, but typically by only 5-10% for any one metal, despite metal geometries that vary from tetrahedral to octahedral, and individual O shift variations with pH of up to 300 ppm. The same principle also works, after due averaging, for related species containing more than one of these metals.

### Oxygen shift sums per metal

Anion Shift s	sum per V	Anion Shift s	sum per Mo	Anion Shift st	ım per W	
VO <sub>4</sub> 3-	2260	MoO <sub>4</sub> <sup>2-</sup>	2116	WO <sub>4</sub> <sup>2</sup> -	1680	
HVO <sub>4</sub> 3-	2292	$Mo_2O_7^{2-}$	2269	$W_6O_{19}^{2-}$	1585	
V <sub>2</sub> O <sub>7</sub> <sup>4</sup> -	2288	Mo <sub>6</sub> O <sub>19</sub> <sup>2-</sup>	2056	W <sub>7</sub> O <sub>24</sub> 6-	1615	
V <sub>3</sub> O <sub>10</sub> 5-	2302	Mo <sub>7</sub> O <sub>24</sub> <sup>2-</sup>	2021	$W_{10}O_{32}^{4}$	1604	
V <sub>10</sub> O <sub>28</sub> 6-	2221	$\alpha$ -Mo <sub>8</sub> O <sub>26</sub> <sup>4</sup> -	2161	$H_2W_{12}O_{42}^{10-}$	1661	
H <sub>2</sub> V <sub>10</sub> O <sub>28</sub> <sup>4</sup> -	2232	β-Mo <sub>8</sub> O <sub>26</sub> <sup>4-</sup>	2090	$\alpha$ -HW <sub>12</sub> O <sub>40</sub> <sup>7</sup> -	1578	
H <sub>3</sub> PV <sub>14</sub> O <sub>42</sub> <sup>6</sup> -	2257*	$PMo_{12}O_{40}^{3}$	2069*	$\alpha$ -H <sub>2</sub> W <sub>12</sub> O <sub>40</sub> <sup>6</sup> -	1583	
H <sub>6</sub> PV <sub>14</sub> O <sub>42</sub> <sup>3</sup> -	2235*	SiMo <sub>12</sub> O <sub>40</sub> <sup>4-</sup>	2063*	PW <sub>12</sub> O <sub>40</sub> <sup>3-</sup>	1605*	
] * 11 72		12 10		SiW <sub>12</sub> O <sub>40</sub> <sup>4-</sup>	1593*	
*ignoring XO <sub>4</sub> unit(s)			$\alpha$ -P <sub>2</sub> W <sub>18</sub> O <sub>62</sub> <sup>6</sup> -	1588*		

The above table provides a theoretical challenge as well as a useful empirical rule. By summing the shifts we are assessing how much they differ from that of water. Very roughly speaking, the shift of any one oxygen is a measure of its covalent bonding to the nearest metal atom. (For this reason it also correlates well, though not linearly, with the reciprocal of the shortest MO bond distance.) The MO bonds have both  $\sigma$  and  $\pi$  components, and show some evidence of delocalisation. Thus there are some analogies with the distribution of (nominally)  $\pi$ -electron spin densities in aromatic systems, where any gain at one carbon is offset by losses at other carbons. One could describe such shifts as resource-limited. Also, the MO bond length sums will remain roughly constant if the metal atoms merely move within an oxygen coordination shell of fixed geometry under the influence of e.g. electrostatic forces. Even so, the shift sum constancies listed above are remarkable, and I would invite readers to offer a deeper explanation of them.

With my best regards

oriver Howarh



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### 01 April 1992 (received 4/6/92)

Dr Bernard Shapiro
TAMU NMR Newsletter
966 Elsinore Court
Palo Alto CA 94303
UNITED STATES OF AMERICA

Dear Dr Shapiro

RE: Lineshape deconvolution - who needs software?

I was in dread of receiving your ultimatum, knowing that I didn't have much material that I was able to talk about. Thankfully I was in the middle of developing a method when your ultimatum arrived.

We are currently producing trimethyl borate (TMB) from boric acid and methanol. The problem has been being able to analyse for TMB and boric acid in the azeotrope by NMR as they have very similar shifts in dimethylacetamide (19.8 and 19.4 ppm respectively). Since we lack deconvolution software we considered exploiting their differences in chemical reactivity. By bubbling ammonia gas through the azeotrope two reactions occur;

 $B(OMe)_3 + NH_3$ 

H<sub>3</sub>N-B(OMe)<sub>3</sub>

19.3 ppm

 $B(OH)_2 + NH_3$ 

 $NH_4^+$  +  $OB(OMe)_3$ 

1.8 ppm

The two species now have resonances well separated (Fig. 1) so that they are readily integrated to obtain ratios of boric acid to TMB in the azeotrope.

Incidentally, if anyone can recommend a good deconvolution package for a 386 or 486 I would be keen to know.

Your sincerely

Roger Meder

Figure 1. Spectrum of TMB and boric acid after ammonia treatment. Shifts are relative to external BF 3.

20 10 PPM

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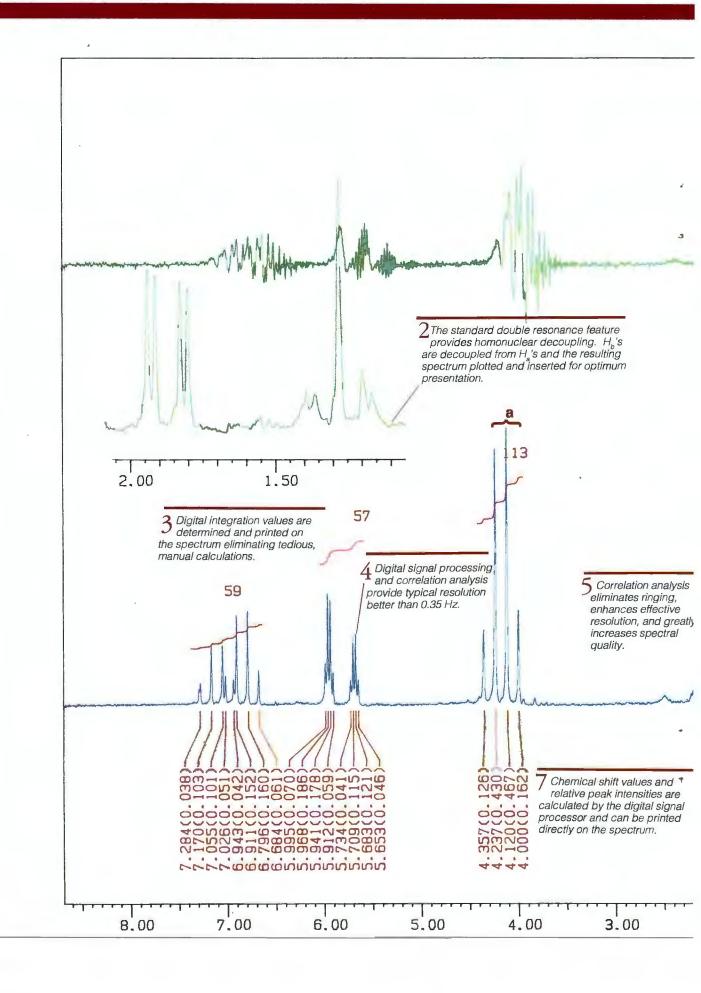
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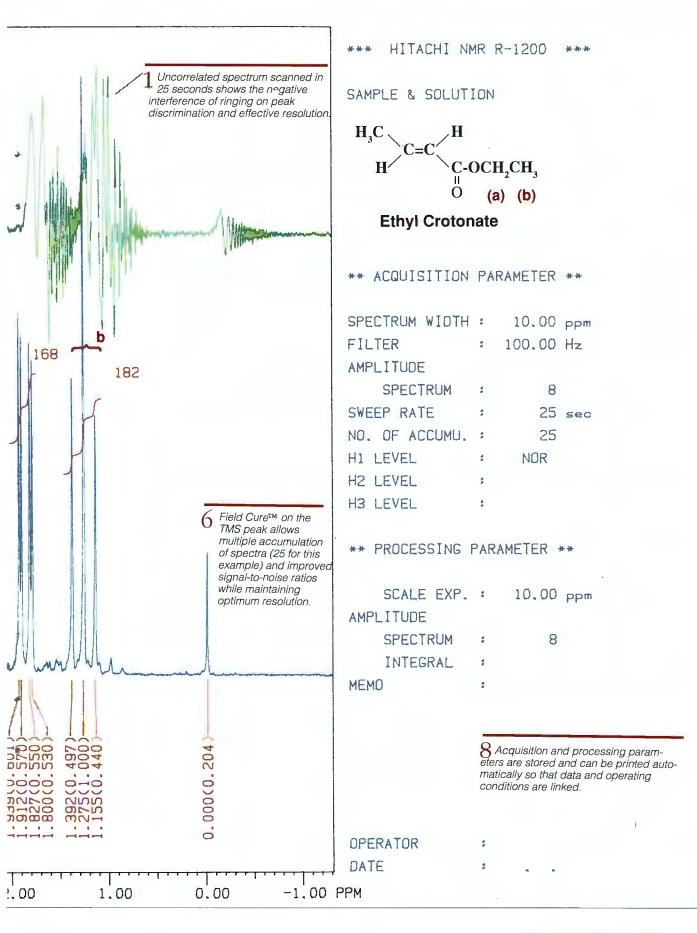
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March 27, 1992 (received 4/3/92)

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Phone +61 3 389 965 FAX +61 3 389 9655

Dr. B.L. Shapiro, TAMU NMR Newsletter, 966 Elsinore Court, Palo Alto CA 94303, USA.

Dear Barry,

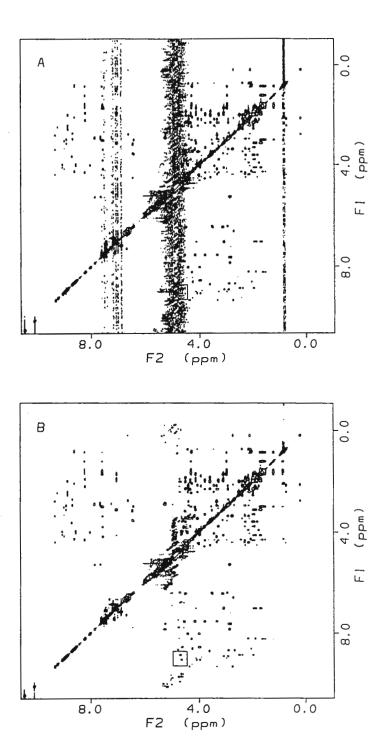
### Suppression of t<sub>1</sub> noise and solvent artifacts in 2D NMR spectra of proteins

Firstly my apologies for provoking the dreaded Utimatum, the threats in which hopefully will be offset by this contribution. One reason for the delayed contribution is my move from the University of New South Wales to the Biomolecular Research Institute in Melbourne. The BRI is a government funded intiative focussing on structural biology and molecular design. The NMR Laboratory is equipped with Bruker AMX-500 and -600 spectrometers and a network of UNIX workstations, supported by a dual processor Convex.

To matters scientific, where I wanted to describe briefly some spectral processing routines written by a PhD student in my group at the University of NSW, Nick Manoleras. The first is a baseline correction routine in which valid baseline segments are identified by their lower slope relative to regions containing peaks. The properties of the these known baseline sections are then used to predict the baseline in areas occupied by peaks or bad t<sub>1</sub> noise. The operation of BFX can be visualized as tracing out a baseline from the data similar to that which an operator would define by following the curve of the baseline with a cursor. It extracts a baseline from the data, then subtracts it to produce a flat baseline. In our experience it gives baseline correction superior to that achievable with polynomial or cubic spline routines.

The other technique is used to reduce t<sub>1</sub> noise artifacts in 2D spectra. To begin with, a t<sub>1</sub> noise profile is produced by measuring the average noise in each column. This profile is then used to determine weighting coefficients for a sliding weighted smoothing filter that is applied to each row, such that the amount of smoothing each point receives is proportional to both its estimated t<sub>1</sub> noise level and the level of t<sub>1</sub> noise of neighbouring points. Thus, points in the worst t<sub>1</sub> noise bands receive the greatest smoothing, whereas points in low noise regions remain relatively unaffected. In addition, weighted smoothing allows points in low noise regions to influence neighbouring points in noisy regions. This method is also effective in reducing the noise artifacts associated with the solvent resonance in spectra of biopolymers in aqueous solution.

The effects of applying the  $t_1$  noise suppression routines to the TOCSY spectrum (recorded on an AM-500) of a polypeptide in  $H_2O$  are illustrated in the Figure. (A) Matrix after baseline correction in  $\omega_2$ . (B) Same as A (and plotted at the same level), but after the application of the  $t_1$  noise suppression routines.



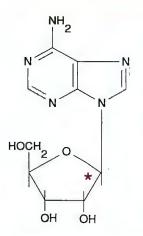
A manuscript describing these procedures has been submitted recently. At present they operate in conjunction with FELIX, but we plan to make them more generally usable.

Yours sincerely,

Raymond S. Norton Head, NMR Laboratory

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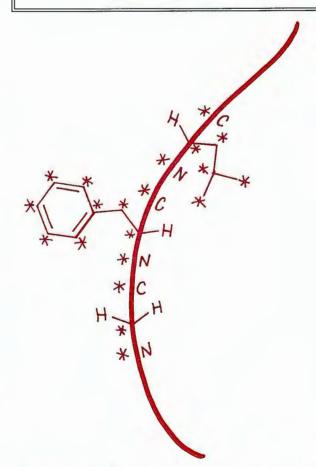
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April 8, 1992

Dr. B. L. Shapiro TAMU Newsletter 966 Elsinore Court (received 4/13/92)

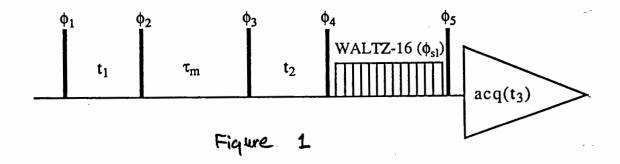
Palo Alto, California 94303

Sensitivity-Enhanced 3D NOESY-TOCSY: Applications to Transferred NOEs

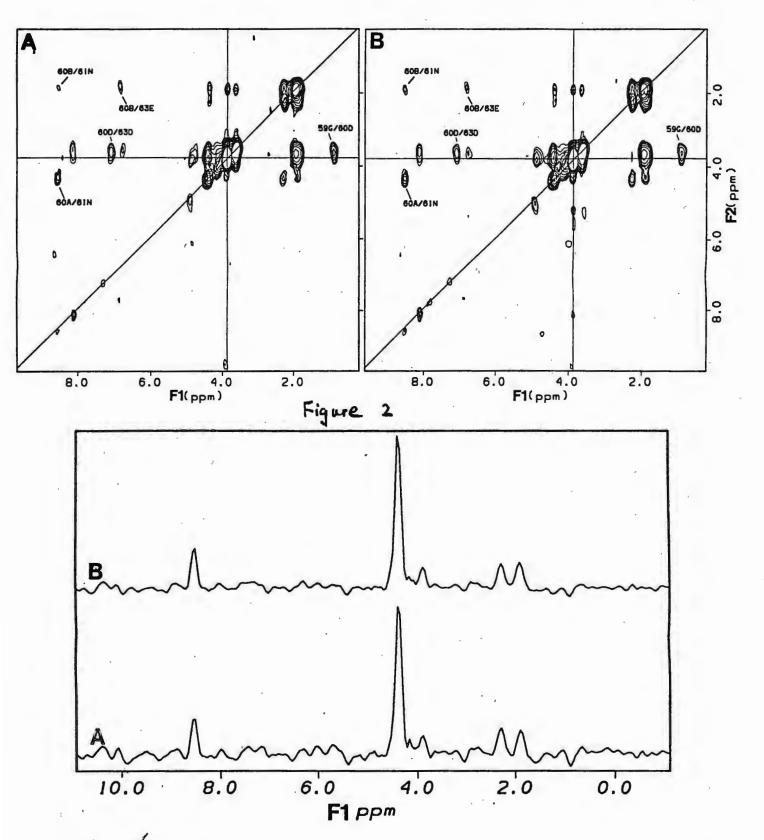
Dear Dr. Shapiro,

In the usual pulse sequence for 3D NOESY-TOCSY, spin-locked coherence transfers are achieved by use of a trimmed MLEV-17 sequence. We replaced the MLEV-17 sequence by a z-filtered WALTZ-16 sequence (Figure 1) to obtain sensitivity enhancement (Ni, JMR 1992, in press). This method involves post-acquisition combinations of the FIDs acquired from subgroups of the full phase cycle utilized for the pulse sequence. The two detected FIDs are stored in seperate memory blocks for subsequent processing. Adequate combinations of these FIDs would restore the two orthogonal components,  $I_{iy}$  and  $I_{ix}$ . If both the  $I_{iy}$  and  $I_{ix}$  FIDs are processed to yield absorptive spectra, they can be combined to produce a spectra with peaks doubled in size compared to each individual spectrum. The key to sensitivity enhancement lies in the fact that noise components in the two spectra are statistically independent and the post-processing combinations are then equivalent to signal averaging, reducing noise in the process (Cavanagh & Rance, JMR & 72, 1990).

Sensitivity-enhanced NOESY-TOCSY was applied to an anticoagulant peptide, acetyl-Asp-Phe-Glu-Glu-Ile-Pro $_{60}$ -Glu $_{61}$ -Glu-Tyr $_{63}$ -Leu-Gln-COOH at a peptide concentration of 6 mM and a thrombin concentration of 0.5 mM in an aqueous solvent of 90%  $H_2O$  and 10%  $D_2O$  at pH 5.5. Figure 2A is the  $F_1$ - $F_2$  plane sliced through the frequency of one of the well-resolved  $\delta CH_2$  protons of Pro along the  $F_3$  dimension of the  $I_{iy}$  3D spectral matrix. The sensitivity-enhanced  $(I_{iy}+I_{ix})$  spectrum is presented in Figure 2B. For some peaks sensitivity ehnancenment is less than the  $\sqrt{2}$  factor expected theoretically. This is probably due to the fact that the  $I_{ix}$  components travel through different transfer pathways and are much more sensitive to prolonged delays both before and after the WALTZ-16 spin-lock pulse (Figure 1).







Feng Ni Ph.D Protein Engineering

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Professor Gerhard Wagner

Tel. (617) 432-3213 or 4366

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240 Longwood Avenue Boston, Massachusetts 02115

April 16, 1992 (received 4/20/92)

Dr. Barry Shapiro TAMU NMR Newsletter 966 Elsinore Court Palo Alto, CA 94303

### Mapping of Spectral Densities of NH Motions in Proteins

### Dear Barry:

We have developed a strategy for direct mapping of spectral density functions from measurements of relaxation parameters of the NH groups in proteins. Previous strategies for interpretation of relaxation parameters were based on mechanistic models of internal motions, such as the "wobbling in a cone" model (1, 2). Alternatively, the "model-free approach" was developed (3,4) which makes, however, the assumption that the autocorrelation function G(t) can be described as a sum of decaying exponentials. As a consequence, the shape of the spectral density function is restricted to a sum of Lorentzians. Such assumptions about mechanistic models or about a functional dependence of the functional dependence of J(w) on certain parameters were necessary since the relaxation rates depend on five values of the spectral density functions, J(0), J( $\omega_N$ ), J( $\omega_H$ + $\omega_N$ ), J( $\omega_H$ + $\omega_N$ ), while fewer experimental <sup>15</sup>N relaxation parameters were measured, such as T<sub>1</sub>, T<sub>2</sub> and the heteronuclear NOE. Thus the problem is underdetermined, and model assumptions have to be made. We added some additional experiments to be able to calculate the values of the spectral density functions (5). We measure six relaxation parameters: the longitudinal relaxation rate of N<sub>Z</sub>, the transverse

underdetermined, and model assumptions have to be made. We added some additional experiments to be able to calculate the values of the spectral density functions (5). We measure six relaxation parameters: the longitudinal relaxation rate of  $N_z$ , the transverse relaxation rates of in-phase and anti-phase coherence,  $N_{x,y}$  and  $2N_{x,y}H_z$ , the relaxation rate of longitudinal two-spin order,  $2N_zH_z$ , the heteronuclear cross relaxation rate and the longitudinal relaxation rates of protons. The pulse sequences to measure these parameters for an  $^{15}N$  enriched protein yield 2D  $^{15}N$ - $^{1}H$  correlated spectra in which the peak intensities can be measured as a function of a relaxation delay  $\tau$ . The analysis includes relaxation by dipole-dipole interaction and CSA relaxation. Crosscorrelation between dipole-dipole and CSA is minimized by appropriate decoupling sequences (6,7). A complete set of data was recorded on a AMX 500 spectrometer for the protein eglin c, and spectral densities were derived (8). The spectral densities decay from  $\omega$ =0 to  $\omega_H$ + $\omega_N$ . For many N-H groups, the spectral density increases significantly from  $\omega_H$ + $\omega_N$  to  $\omega_H$ - $\omega_N$  (due to the negative sign of  $\gamma_N$ , the latter is the larger frequency). Should this unexpected observation prove to be real, our common picture of how an autocorrelation function should look like had to be revised.

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Sincerely

Joseph W. Pers Jeffrey W. Peng Gerhard Wagner

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All Newsletter Correspondence

Should Be Addressed To:

Dr. Bernard L. Shapiro TAMU NMR Newsletter 966 Elsinore Court Palo Alto, CA 94303, U.S.A.

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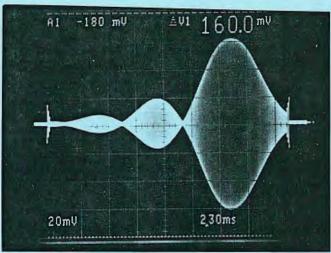
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# Gradient Enhanced Spectroscopy

### Phase Sensitive DOF COSY

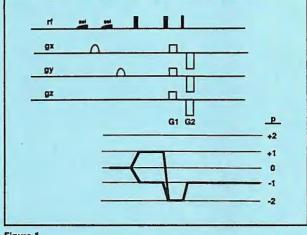
The selection of multiple quantum coherence with gradients is an effective method for suppressing the water resonance in aqueous solutions and for reducing t<sub>1</sub> noise and other artifacts. A further enhancement is the use of selective water excitation using crafted RF pulses followed by gradient dephasing to attenuate the water signal prior to the coherence selection sequence. This technique has been applied to the phase sensitive gradient enhanced DQF COSY experiment for a 5mM lysozyme sample in 90% H<sub>2</sub>O. The data were collected on an Omega PSG 500 equipped with the S-17 Gradient Enhanced Spectroscopy accessory using a 5mm inverse probe.



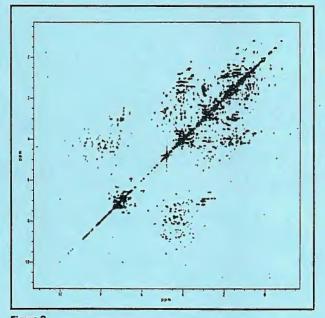
The crafted RF pulse used in this sequence. This pulse was designed using the "hard pulse approximation" method. It is characterized by a flat amplitude response in the selected region and minimal excitation in the out of band region. The duration of the selective RF pulse was set to 20 ms, corresponding to an excitation bandwidth of 175 Hz.

Figure 3

A phase sensitive GE-DQF COSY spectrum of 5 mM lysozyme in 90% H<sub>2</sub>O. Thirty-two scans were accumulated for each of the 700 t<sub>1</sub> increments resulting in a total data acquisition time of approximately 9.7 hours. Half-sinusoidal gradient pulses of 20 ms duration with an amplitude of 10 G/cm were used to dephase the excited water signal. Coherence selection was achieved using 2 ms gradient pulses of 17



The pulse sequence and the corresponding coherence level diagram for the phase sensitive GE-DQF COSY experiment. The water resonance is selectively excited using a crafted RF pulse followed by a shaped gradient pulse along the X axis which dephases the transverse magnetization. A second selective RF pulse is applied with a dephasing gradient along the Y axis to evoid gradient recalled echoes. This is followed by a phase sensitive GE-DQF COSY sequence.

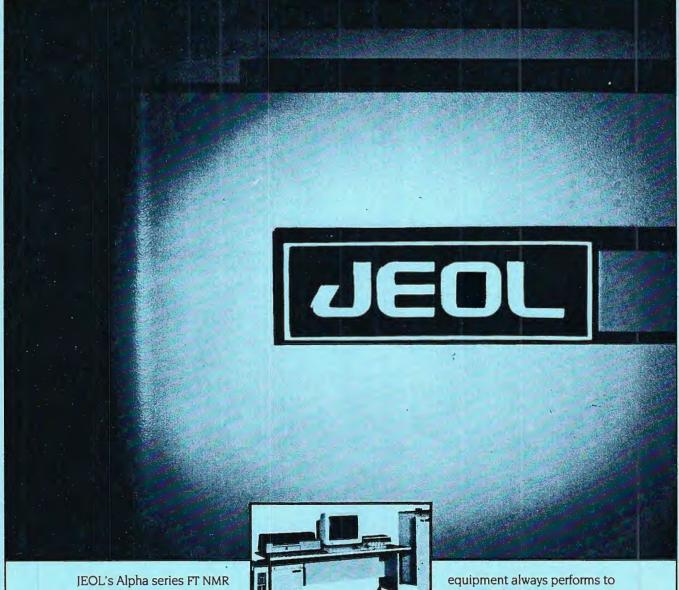


and 34 G/cm amplitude.



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