

Jim Lippmaa
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Newsletter

No. 204

September, 1975

- B. L. Shapiro
Policies and Practical Considerations. . . 2
- H. Pearson
Downfield δ Shifts in ^{13}C NMR of Ortho-
substituted Toluenes 4
- E. Lippmaa and M. Alla
High Resolution Broad Line ^{13}C NMR in
Solids 5
- D. Ruben
NMR Powder Study of Molecular Motion in
a Six-Fold Potential 7
- A. Allerhand
Effect of ^{13}C - ^{14}N Dipolar Interactions
on T_1 Values and Intensities of Non-
protonated Carbon Resonances. 9
- N. Cyr, G. Ritchie and A. S. Perlin
Assignment of ^{13}C Signals Using a ^1H
Coupled Spectrum. 12
- N. V. Riggs
A Trinuclear Varian HA60-IL Hybrid. . . . 15
- I.K. O'Neill, H.J. Prosser and C.P. Richards
Quantitative C-13 FT-NMR?. 17
- P. C. Lauterbur and L. J. Altman
A Tritium NMR Facility. 19
- K. G. Sharp and K. L. Servis
 ^{29}Si - $\{^{19}\text{F}\}$ NOE and T_1 for Hexafluorodi-
silane 22
- P. C. Lauterbur
Zeugmatographic Spectroscopy. 23
- H. J. C. Yeh
NMR Study of the Mechanism of Isomer-
ization of a Series of Aldimines 25
- J. D. Roberts
Natural-Abundance ^{15}N NMR with the WH-180. . 27
- G. N. LaMar
Shift-Reagent-Induced Relaxation. . . . 29
- M. D. Johnston, Jr. and D. J. Raber
Structure Elucidation with Lanthanide-
Induced Shifts. 31
- E. B. Whipple
Random Noise Stirring 34
- F. Hasan, R. Deslauriers and I.C.P. Smith
A Convenient Means to Remove the
Influence of Residual Paramagnetic Metal
Ions in Deuterium Oxide on ^{13}C Spin-
Lattice Relaxation Times 35
- J. T. Gerig
Surplus Tubes for Sale 38
- R. Wasylshen
 ^1H - T_1 's in Condensed Ring Hydrocarbons. . 39
- E. Martinelli and A. Ripamonti
 ^{13}C Spin Lattice Relaxation Times of
the Antibiotic Rifampin. 41
- S. B. W. Roeder and T. P. Higgs
Computationally Offsetting the Frequency
in NMR Spectra and Other NMR Signals . . 45
- R. K. Harris and R. H. Newman
 ^{13}C Spin-Lattice Relaxation for 1-Phenyl-
adamantane 47
- D. N. Lincoln and V. Wray
Measuring Accurate Line Positions on CFT-20. 49
- R. H. Cox
N-Nitroso-N-methylaniline, An Error. . . 51
- G. E. Wilson, Jr.
Rotameric Preference of β -Alanine . . . 52

Bruce Hawkins ✓

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Deadline Dates: No. 205: 6 October 1975

No. 206: 3 November 1975

All Newsletter Correspondence, Etc. Should Be Addressed To:

Dr. Bernard L. Shapiro
 Department of Chemistry
 Texas A&M University
 College Station, TX 77843 U.S.A.

AUTHOR INDEX - TAMU NMR NEWSLETTER NO. 204

Alla, M.	5	Perlin, A. S.	12
Allerhand, A.	9	Prosser, H. J.	17
Altman, L. J.	19	Raber, D. J.	31
Cox, R. H.	51	Richards, C. P.	17
Cyr, N.	12	Riggs, N. V.	15
Deslauriers, R.	35	Ripamonti, A.	41
Geric, J. T.	38	Ritchie, G.	12
Harris, R. K.	47	Roberts, J. D.	27
Hasan, F.	35	Roeder, S. B. W.	44
Higgs, T. P.	44	Ruben, D.	7
Johnston, M. D.	31	Servis, K. L.	22
LaMar, G. N.	29	Shapiro, B. L.	2
Lauterbur, P. C.	19 & 23	Sharp, K. G.	22
Lincoln, D. N.	49	Smith, I. C. P.	35
Lippmaa, E.	5	Wasylishen, R.	39
Martinelli, E.	41	Whipple, E. B.	34
Newman, R. H.	47	Wilson, G. E.	52
O'Neill, I. K.	17	Wray, V.	49
Pearson, H.	4	Yeh, H. J. C.	25



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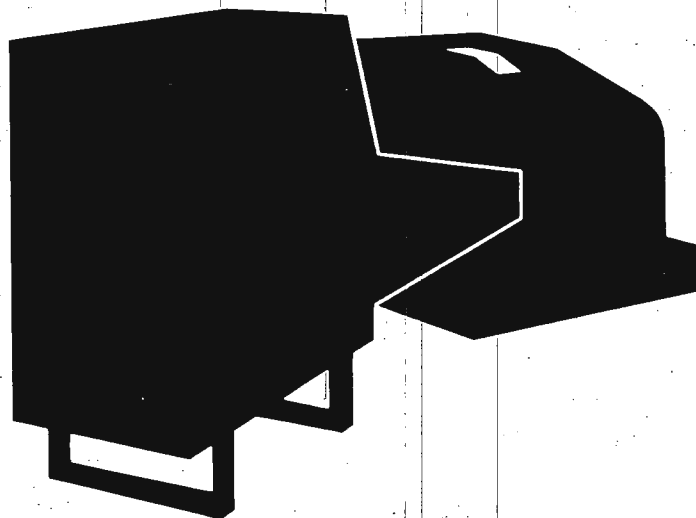
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Participation is the prime requisite for receiving the TAMU NMR Newsletter; in order to receive the Newsletter, you must make at least occasional contributions to its contents. We feel that we have to be ruthless in this connection and the following schedule is in effect: Eight months after your last contribution you will receive a "Reminder" letter. If no contribution is then forthcoming ten months after your last contribution, you will receive the "Ultimatum" letter, and then the next issue will be your last. If you are dropped from the mailing list, you can be reinstated by sending a contribution, and you will receive back issues (as available) and forthcoming issues at the rate of nine per contribution. Frequent contributions are encouraged, but no "advance credit" can be obtained for these. In cases of joint authorship, either contributor, but not both, may be credited - please indicate to whose account credit should be given.

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4. Practical Considerations: (a) All contributions to the TAMU NMR Newsletter should be sent to the undersigned and will always be included in the next issue if received before the deadline dates, which appear in each issue.

(b) Contributions should on the minimum (NOTE!!!) number of 8½ x 11" (21 x 27.5 cm) pages printed on one side only. Margins should be between 2 and 3 cm on all sides - PLEASE observe these limits if at all possible. Black ink, typing, drawings, etc., essential. We are not equipped to deal with large size pieces of paper - e.g., A-60 charts.

Please conserve space by avoiding double spacing (except where necessary), ultra-wide margins, half-filled pages, etc. In general, please plan and construct your contribution so as to fill the minimum number of pages needed. On the other hand, drawings and spectra lose both eye-appeal and utility when they are too small. Only in very rare and absolutely necessary circumstances will a contribution in excess of three pages - including drawings, figures and references - be accepted. Economic necessity forces this policy.

Since reproductions of various kinds do not themselves reproduce too well, contributors are urged to submit their photographic originals to us (if the size does not exceed 8½ x 11"), and we will be happy to return these if requested. Some law of physics says that photographic reproductions of fuzzy or blurred originals never come out less fuzzy or blurred.

(c) Please provide short titles of all topics of your contributions, as they will ensure accuracy in preparing the title-page index.

(d) Please do not send in manuscripts, theses, books, etc., and ask us to be your consciences in selecting what should and shouldn't go into the Newsletter.

5. Suggestions: They are always welcome.

B. L. Shapiro

B. L. Shapiro
1 August 1975

Address for all contributions and inquiries:

Professor Bernard L. Shapiro
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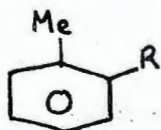
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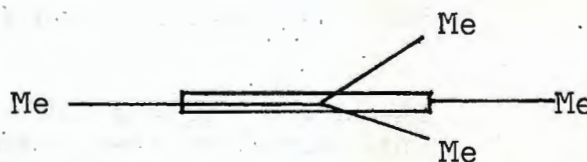
Downfield δ Shifts in ^{13}C NMR of Orthosubstituted
toluenes.

Downfield 'steric' shifts have attracted some interest recently¹. The shift of the toluene methyl carbon on ortho alkyl substitution is interesting in this respect and some values are listed below.

$\delta(\text{Me})$, ppm from TMS



R=H	21.4
R=Me	19.4
R=Et	19.0
R=iPr	19.2
R=tBu	23.2



Substitution in toluene of an ortho methyl group produces the well known upfield shift of 2.0 ppm; however, substitution of ortho ethyls or ortho isopropyl (that is addition of 1 or 2 δ methyl groups) produces an almost identical shift. Thus the δ effects of the added methyls are near zero. In contrast the addition of a third δ methyl, as in the t-butyl compound, produces a large downfield shift of 3.8 ppm, a magnitude not dissimilar to those noted previously^{1,2,3} where syn axial interactions were involved. It is unlikely that a pure syn axial interaction is involved here, the shift probably reflects the consequences of two quasi syn axial interactions in a conformation such as 1. On the other hand the data suggest that the dominant conformation in the isopropyl compound lacks any 1,5 Me-Me interactions.

Thus the generality of δ interactions is further established and in this type of system they are a useful indicator of conformation and caution is again drawn to the fact that it is dangerous to associate steric crowding with upfield shifts.

Please credit this to Ray Abraham's account.

1. J.G. Batchelor, J. Magn. Resonance 18, 212 (1975). With all good wishes
2. S.H. Grover et al, J. Magn. Resonance 10, 227 (1973). Harry Pearson.
3. J.I. Kroschwitz et al, J. Amer. Chem. Soc. 91, 5927 (1969).

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№ 44

1st August 1975

Prof. Bernard L. Shapiro
Department of Chemistry
Texas A & M University
College Station,
Texas 77843
USA

Dear Professor Shapiro,

High Resolution Broad Line ^{13}C NMR in Solids

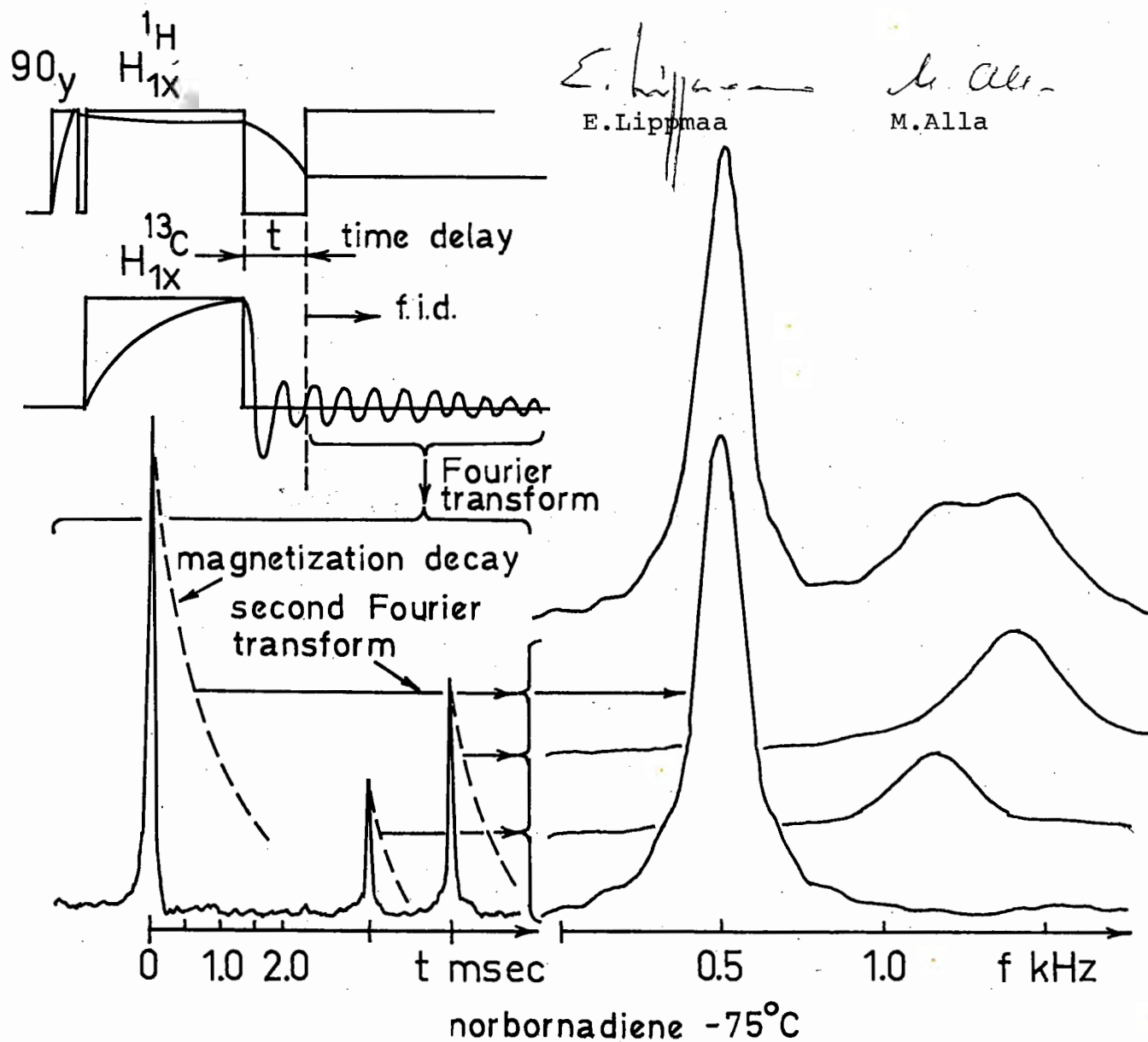
Thank you for the blue reminder. We have recently become interested in the proton-enhanced nuclear induction spectroscopy, developed by J.S. Waugh and A. Pines. This new tool already combines high sensitivity with high selectivity, thus having the best of both, in double resonance at least. It seemed like a good idea to extend these advantages to single resonance as well, in order to study lineshapes, relaxation and other aspects of solid-state dynamics. To do this, the abundant proton spins are spin-locked in the rotating frame in the usual manner. A single mutual contact with the rare (^{13}C) spins is established by a strong rf field, fulfilling the Hartmann-Hahn condition, but the free induction decays are co-added and Fourier transformed in a NIC-1085, NIC-293 computer only after a variable delay, during which the decoupling field is also turned off temporarily. The free induction decays are registered in the decoupled mode, thus yielding a series (12 to 20) high resolution solid state spectra. The (delay) time dependences of line intensities in these spectra represents the transverse relaxation of groups of nonequivalent ^{13}C nuclei. Another Fourier transform of data points, corresponding to these time dependences, gives the shapes of all individual uncoupled lines that can be resolved in the decoupled spectrum. The whole "single resonance" spectrum can thus be reconstituted, but with about 100-fold increased sensitivity and high selectivity.

A low temperature study of solid norbornadiene, which is a semispherical molecule with fast rotation even at low temperatures, showed that the solid-state ^{13}C lines have different widths with different temperature dependences and shapes that are neither Lorentzian nor Gaussian. The polarization transfer from ^1H to ^{13}C nuclei is markedly non-exponential and characterized by different time constants for all groups of nonequivalent carbons in the molecule.

A more full account is going to appear in Chemical Physics Letters.

With best regards,

Sincerely yours,



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DEPARTMENT OF CHEMISTRY

BERKELEY, CALIFORNIA 94720

August 4, 1975

Professor B. L. Shapiro
Department of Chemistry
Texas A & M University
College Station, Texas 77843

Dear Professor Shapiro:

RE: NMR Powder Study of Molecular Motion in a Six-Fold Potential

We have recently joined some other groups in examining the details of molecular reorientation in solids by observing the effect of the motion on the powder lineshapes. The theory has been discussed by Luz et al.¹ (LBA), and by Spiess.² The experimental technique is ¹³C nmr using proton-enhanced nuclear induction spectroscopy.

Hexamethyl benzene was one of the first compounds we studied. Motion in the solid is about the 6-fold axis. In the zero motion limit a typical asymmetric powder pattern should be seen; in the fast motion extreme the pattern should be that of a symmetric shielding tensor. The interesting range is when the molecules are just beginning to reorient. In this case different patterns will be seen depending upon whether the molecules rotate diffusively in the plane of the ring or make rapid jumps of 60° between the equivalent sites. In the latter case, as predicted by LBA,¹ two anomalous "bumps" should be seen in the powder pattern between σ_{11} and σ_{22} . These (fondly referred to in one laboratory as "shpitzim") result from molecules oriented so that two of the three different ring carbons have approximately the same chemical shift. For these molecules even very low jump frequencies will satisfy the exchange narrowing condition.

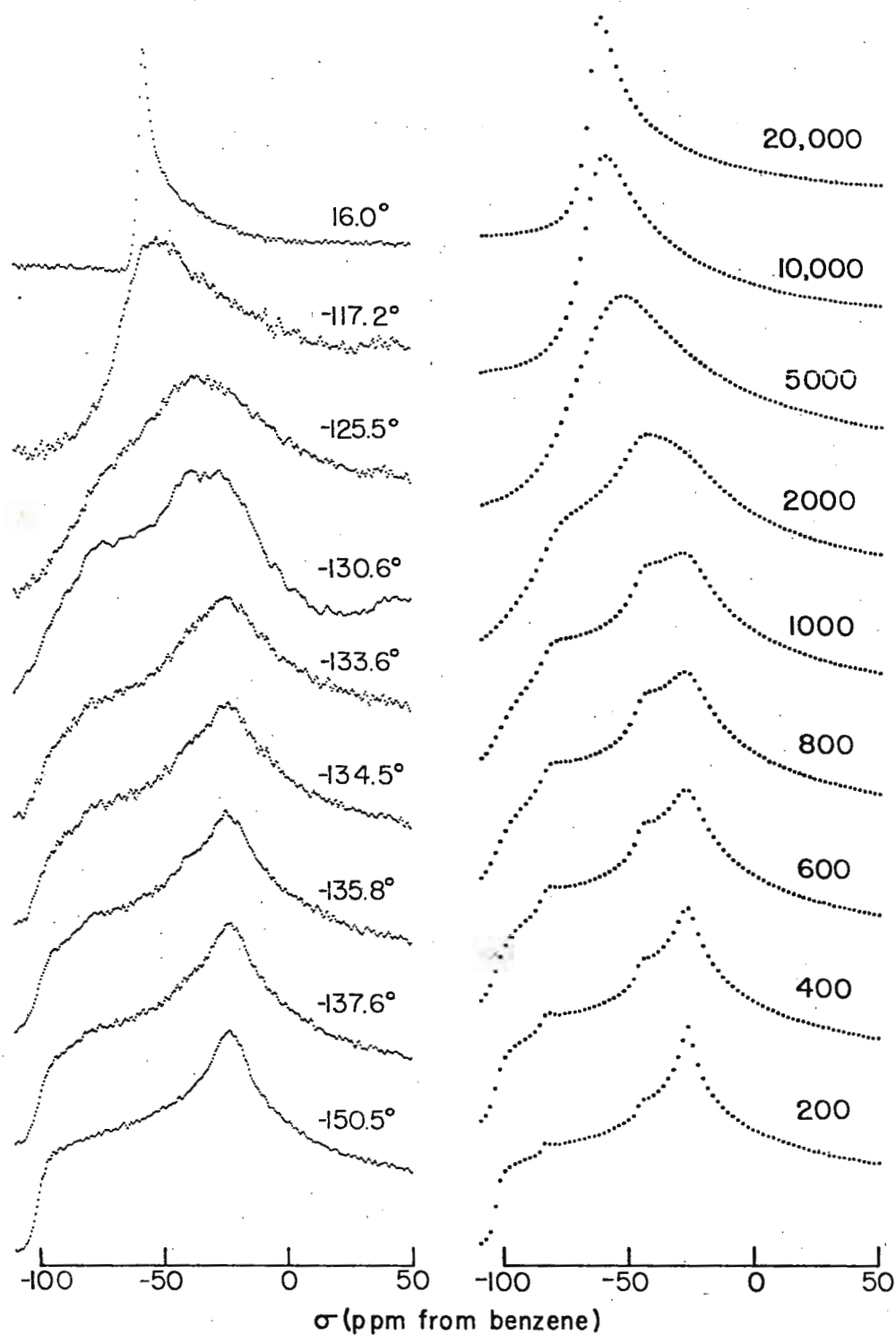
Some of our results are shown in the accompanying figure. The left side shows experimental spectra at various temperatures and the right side theoretical lineshapes at various jump frequencies in Hz. These we calculated using the standard exchange matrix formalism with powder orientation averaging by computer, and agreement is seen to be good. The spectra are also in agreement with those calculated by the elegant expansion techniques of LBA.¹ The bump at ~ -80 ppm is especially evident in the -130.6° spectrum. Lineshapes calculated assuming a rotational diffusion model were in total disagreement with experiment.

Most of the experimental work was done by David Wemmer, a graduate student in our laboratory. Our results on the motion in HMB and other molecules will be published shortly.

Best regards,

Dr. David Ruben

1. Z. Luz, A. Baram and S. Alexander, should be appearing about this time.
2. H. W. Spiess, Chemical Physics 6, 217 (1974).



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Please credit this to the account of Professor Alex Pines

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TEL. NO. 812—

August 6, 1975

Professor Bernard L. Shapiro
Department of Chemistry
Texas A & M University
College Station, Texas 77843

Effect of ^{13}C - ^{14}N Dipolar Interactions
on T_1 Values and Intensities
of Nonprotonated Carbon Resonances

Dear Barry:

We have shown (R. S. Norton and A. Allerhand, JACS, in press) that ^{13}C - ^{14}N dipolar interactions can contribute significantly to the relaxation of a nonprotonated carbon that is directly bonded to one or more nitrogen atoms. We have shown that if the T_1 values of the protonated carbons of a rigid large organic molecule are measured, then one can make fairly accurate predictions of the T_1 and NOE values of nonprotonated carbons within the same molecule, without invoking relaxation mechanisms other than the ^{13}C - ^1H and ^{13}C - ^{14}N dipolar ones (at 14.2 kG). As far as we know, earlier workers have not considered the possible importance of ^{13}C - ^{14}N dipolar contributions to ^{13}C relaxation (of nonprotonated nitrogen-bearing carbons), probably because of the very low gyromagnetic ratio of ^{14}N . However, the short C-N bond length partly compensates for the low gyromagnetic ratio.

Spin-lattice relaxation times and integrated intensities of the resonances in proton-decoupled natural-abundance ^{13}C Fourier transform nmr spectra of adenosine-5'-monophosphate and guanosine-5'-monophosphate (in H_2O and D_2O , at 40-44°, at 15.18 MHz, in 20-mm sample tubes) were compared with calculated values that take into account ^{13}C - ^1H and ^{13}C - ^{14}N dipolar relaxation. In each case, T_1 values of methine carbons of the base were used to obtain a rotational correlation time, which

was then used, together with interatomic distances from crystallographic data, to compute T_1 values of nonprotonated carbons. Nonprotonated carbons which are directly-bonded to nitrogens and which have no hydrogens two bonds removed yielded theoretical T_1 values strongly affected by ^{13}C - ^{14}N dipolar interactions. For carbons in this category, calculated T_1 values which include ^{13}C - ^{14}N dipolar interactions are in much better agreement with experimental values than calculated values which consider only ^{13}C - ^1H interactions. Integrated intensities were calculated by considering variations in the nuclear Overhauser enhancement that result from differences in relative contributions to $1/T_1$ from ^{13}C - ^1H and ^{13}C - ^{14}N dipolar relaxation. The calculated intensities are in excellent agreement with the experimental ones.

Best regards,



Adam Allerhand
Professor of Chemistry

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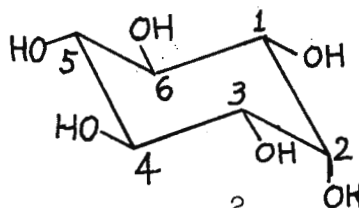
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Dr. Barry Shapiro,
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Dear Dr. Shapiro:

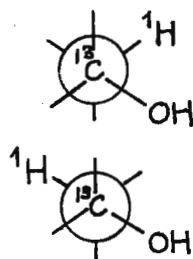
Assignment of ^{13}C Signals Using a ^1H Coupled Spectrum

We would like to demonstrate that two-bond and three-bond ^{13}C -H coupling constants ($^2J_{\text{CH}}$ and $^3J_{\text{CH}}$, respectively) can be used to assign the ^{13}C signals unambiguously. When all the protons are decoupled, myo inositol has four ^{13}C signals with the intensity ratios 1:2:1:2 from the low field.



Knowing the following², one can simply predict the ^{13}C signal patterns when protons are coupled:

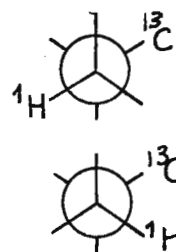
$^2J_{\text{CH}}$



$\sim 5 \text{ Hz}$

$\sim 0 \text{ Hz}$

$^3J_{\text{CH}}$



$\sim 5 \text{ Hz}$

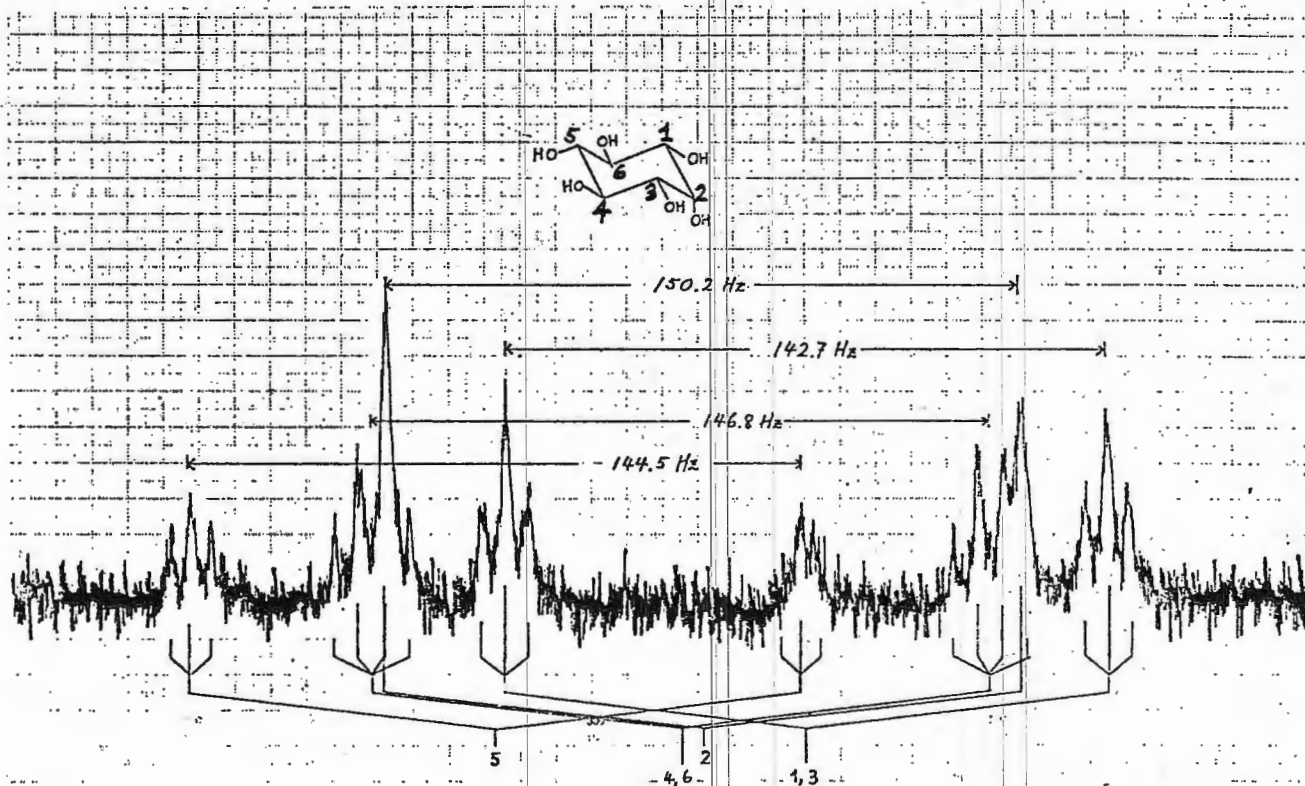
$\sim 0 \text{ Hz}$

Dr. Barry Shapiro

August 6, 1975

Carbon	Expected Couplings	Doublet Patterns*
2	$^1J_{CH(2)}$	singlet
5	$^1J_{CH(5)}, ^2J_{CH(4,6)}$	triplet
1	$^1J_{CH(1)}, ^2J_{CH(2,6)}$	triplet
4	$^1J_{CH(4)}, ^2J_{CH(3,5)}, ^3J_{CH(2)}$	quartet

* Each carbon resonance will exhibit a large doublet due to one directly bonded proton.



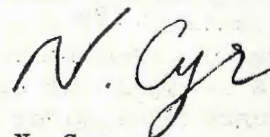
The 1H coupled ^{13}C spectrum with assignments is shown above. The chemical shift results agree with the general observation,^{2,3} that the carbon nuclei next to axial OH groups (C1,3) appear at higher field than those next

Dr. Barry ShapiroAugust 6, 1975

to equatorial ones, and that cis arrangements of vicinal OH groups also shift the carbon resonances involved to high field. This assignment for C1,3 and C4,6 is the reverse of that given in ref. 1.

1. D.E. Dorman, S.J. Angyl and J.D. Roberts, J. Amer. Chem. Soc., 92, 1351(1970).
2. N. Cyr, R.G.S. Ritchie, N.K. Richtmyer and A.S. Perlin, 58th C.I.C. Meeting, Toronto, May 1975.
3. A.S. Perlin, B. Casu and H.J. Koch, Can. J. Chem., 48, 2596(1970).

Yours sincerely,



N. Cyr
G. Ritchie
A.S. Perlin

/ce





NVR.RC

THE UNIVERSITY OF NEW ENGLAND

ARMIDALE, N.S.W.

DEPARTMENT OF ORGANIC CHEMISTRY.

7th August, 1975.

Professor B.L. Shapiro,
Department of Chemistry,
Texas A&M University,
COLLEGE STATION,
Texas, 77843, U.S.A.

Dear Professor Shapiro,

Title: A Trinuclear Varian HA60-IL Hybrid
Subtitle: Teaching an Old Dog New Ways to Perform His Tricks
or
Resuscitation for Machines (or Persons) in extremis

We have already described¹ how our aging (born, 1966) Varian HA60-IL spectrometer was converted for ^{13}C n.m.r. measurements in c.w. mode at 15.09 MHz. The appropriate proton-decoupling frequency near 60.0 MHz was generated by mixing a variable frequency near 3.6 MHz with a crystal-controlled frequency near 56.4 MHz, the ^{19}F resonance frequency at the designated field (14092G) for the ^{13}C and ^1H resonance frequencies mentioned. This method was chosen because we proposed to instal an external ^{19}F field-frequency lock. The frequency source for the latter must however be coherent with the ^{13}C exciting frequency and, in these days of financial stringency, we must achieve coherency through use of a General Radio Model 1170 Frequency Synthesizer (FS 1170) (obtained in a relatively affluent yesteryear) and such devices as we can construct with the much smaller funds now available to us.

The FS 1170 provides one dial-selectable and several fixed output frequencies all derived coherently from the 5 MHz internal (or, if desired, external) master crystal. We chose, as the simplest course, to multiply the fixed 1 MHz output to 14.000 MHz (actually in two stages, X7 followed by a doubler) which, in order to avoid possible complications of field modulation on ^{13}C spectra measured in PFT mode, is phase modulated (by a stable a.f. in the range, 1-3 kHz; the manual oscillator of the standard "lock box" is adequate), passed through a limiting amplifier (to become in effect frequency modulated), and fed to the (retuned) V4311 r.f. unit in place of the standard 14.111 MHz crystal source for ^{19}F resonance; the V4311 standard i.f. of 5 MHz is generated by using a 51.000 MHz crystal in the local oscillator. With use of the standard lock-box circuits, the field (ca 13980G) now locks strongly to the ^{19}F resonance of C_6F_6 in a non-spinning 5mm sample tube. The appropriate proton-decoupling carrier frequency (near 59.5 MHz) may now be obtained by mixing a variable frequency near 3.5 MHz with the 56.0 MHz. The corresponding ^{13}C exciting frequency (near 14.97 MHz) is obtained by dialing one-half the required value on the front panel of the FS 1170 and feeding the output to our pulse system (where the frequency is doubled) but, because of delays in completing construction, we have not yet recorded any ^{13}C spectra in PFT mode.

Alternatively, by changing a few BNC connections we can apply field modulation in both analytical and lock channels, and use the standard V4333 5mm probe for $^1\text{H}/^{19}\text{F}$ work to record ^{19}F spectra (not proton-decoupled) in homopolar-internally locked c.w. mode at 56.000 MHz; suitably high audiofrequencies to avoid

Professor B.L. Shapiro, Texas A&M University.

7th August, 1975.

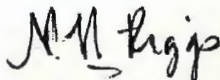
overlap of chosen sideband and other-band spectra, and the convenience of computer control, are provided by the system (apart from the probe) and program used¹ for ^{13}C spectra in c.w. mode.

We now possess an EM-360 for routine, ambient-temperature p.m.r. spectra, but at times shall have to use the (modified) HA60-IL system for higher-resolution and/or variable-temperature p.m.r. spectra. We abhor retuning the V4311 r.f. unit (several retunings have proved deleterious to the tuning slugs) and cannot afford a new, separate 60 MHz unit (if they are still available!). Our approach has been to construct a module containing a 4 MHz crystal-controlled oscillator whose output is fed to two low-noise mixer circuits (6AK5 tubes, as used in the probe preamplifier), one between transmitter and probe, the other between probe and receiver, each with appropriately tuned output stages and levels being matched to those in the original (unmixed) system by tuned low-noise amplifiers. The module is connected or disconnected by changing four BNC connections. The mixer naturally introduces additional noise into p.m.r. signals, but adequate tests of the complete system have been prevented by recent breakdowns in other components. In retrospect, additional noise would probably be more readily tolerable in the ^{19}F lock signal or (infrequently required) ^{19}F spectra, and we intend to retune the system so that the V4311 unit operates at a synthesizer-based frequency of 60.000 MHz and the 4 MHz mix-in frequency for ^{19}F lock must of course be coherent with it.

The system is unquestionably a hybrid and it is at least trinuclear; other nuclei could be added by appropriate trickery, but you will agree that we are approaching, if not already in, extremis.

Dr. D.M. Doddrell was heavily involved in the earlier design stages and Dr. A. Moritz gave valuable advice on phase modulation. Mr. F.B. Hanson and Mr. R.J. Kenny have designed some of the circuits and done all the construction work.

Yours sincerely,


N.V. RIGGS.

¹ Doddrell, D.M., Hanson, F.B., Kenny, R.J., Marker, A., and Riggs, N.V., Aust.J.Chem., 1974, 27, 2175-90.



Department of Trade and Industry
Laboratory of the Government Chemist
 Cornwall House Stamford Street London SE1 9NQ

Telephone 01-928 7900 ext 649

Dr Bernard L Shapiro
 Dept. of Chemistry
 Texas A and M University
 College Station
 TEXAS 77843 USA

Your reference

Our reference

Date 11 August 1975

Dear Dr Shapiro

QUANTITATIVE C-13 FT-NMR?

Recently we have been studying the application of FT-NMR to quantitative analysis. One possible area of application is the analysis of sugar mixtures. As yet we have only studied aqueous sucrose in detail, and have been able to obtain results within 3.8% of the true values for concentrations in the range 0.34 - 1.47 M using 1,2-dihydroxybenzene as a standard. The results are summarised in Table 1. The peak area for C-3, C-6, at 117.3 ppm and C-4, C-5 at 122.1 ppm of 1,2-dihydroxybenzene and C-2' at 104.4 ppm of sucrose were measured (Chemical shifts in ppm downfield from TMS).

Our procedure has been to make up the solutions in 0.1 M potassium hexacyanochromate in order both to eliminate NOE's and to reduce the T_1 's of the nuclei being studied to less than 0.3 sec. at 32°C. The sucrose concentration changes the solutions viscosity markedly and affects the T_1 values as would be expected from the dependence of T_1 on the rate of molecular motion. Measurements have been made using 90 pulses with pulse delay times of 7-10 times the longest T_1 's. Gated decoupling was used to eliminate NOE's not removed by the presence of the relaxation agent.

By chance, in each solution that we examined, the T_1 's for all nuclei specified above were approximately the same. However we have also worked with maltose using the same concentrations and procedure, but where the T_1 for maltose C-1 at 101.1 ppm was approximately half the value for the carbon nuclei of 1,2-dihydroxybenzene.

Initial measurements indicate that the determinations are only within 10-15% of the true concentrations. The integrated peak area ratios for C-1 (maltose) : C-3 + C-6 (1,2-dihydroxybenzene) were found to vary with acquisition time and also the exponential filter used for apodization; the peak areas of the nuclei with long T_1 's increase with increasing acquisition time and smaller apodization functions. It appears that for quantitative determinations by C-13 FT-NMR, it may be necessary for the T_1 's of the nuclei being studied to be essentially equal, or for the data acquisition process to be taken into account. These conditions did not appear to be necessary for quantitative P-31 FT-NMR determinations of phosphorylcolamine ($H_2NCH_2CH_2OPO_3H$) using

$P_{27}^{O^{4-}}$ as an internal standard.

Our experience has been that NOE is not a difficult problem in FT-NMR quantitation (note that we compare protonated and non-protonated carbons), but is lengthy because of the necessary gated decoupling. We note that there is a dearth of quantitative as opposed to semi-quantitative work in the literature and would like to know of other people's experience in this area.

We acknowledge the Government Chemist for authorisation to release these details. Please credit this communication to I K O'Neill.

Yours sincerely

I. K. O'Neill

I K O'NEILL

H. J. Prosser

H J PROSSER

C. P. Richards

C P RICHARDS

Table 1

Quantitative determination of sucrose by C-13 FT-NMR

Sucrose in solution	Sucrose determined	% error
1.466 M	1.519 M	+ 3.6
1.231 M	1.201 M	- 2.4
0.953 M	0.954 M	+ 0.10
0.706 M	0.713 M	+ 1.0
0.344 M	0.357 M	+ 3.8
0.480 M	0.477 M	- 0.6

Stony Brook

August 12, 1975

Professor B. L. Shapiro
Department of Chemistry
Texas A & M University
College Station, Texas 77843

Dear Barry:

"A Tritium NMR Facility"

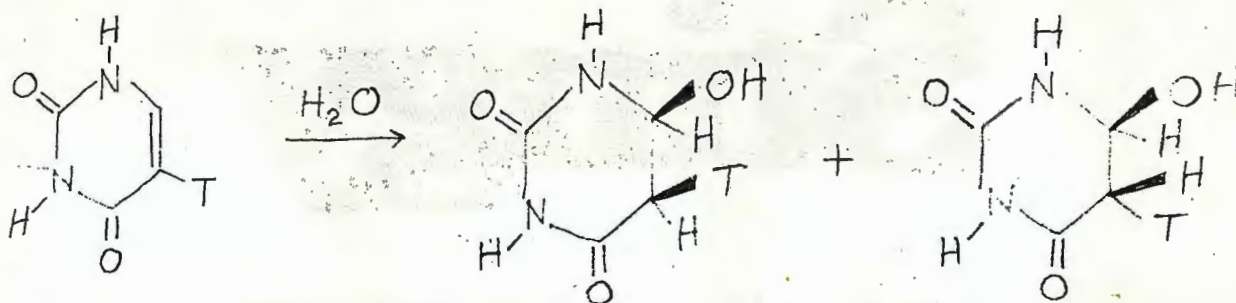
We are in the early stages of setting up a tritium NMR facility centered around a Varian XL-100-12 FT spectrometer and an adjacent hot lab. Our goals include the study of biochemical mechanisms and biosynthetic pathways and the investigation of very complex biological systems in which the high sensitivity and freedom from background signals of ^3H NMR spectroscopy may make possible experiments that cannot be done with other nuclei. The unusual and highly specialized nature of such a laboratory suggests to us that it should be made widely available to the scientific community, and that its operation might be partially funded by a national agency as a special research resource.

Suggestions, comments, and project proposals are hereby solicited. The pace and scope of our planning will depend upon the amount of use that might be made of such a laboratory, and upon the kinds of experiments proposed and the requirements they place upon the spectrometer and upon other laboratory apparatus and facilities. For those of you who have not previously considered the properties of tritium resonances, the main features may be summed up rather simply. The triton behaves like a better proton (except for its instability - a half-life of 12.5 years). The spin is $1/2$ and the magnetic moment slightly larger than that of the proton, so that our XL-100 sees tritium at nearly 107 MHz. One can therefore expect a slightly better sensitivity for ^3H than for ^1H signals, with chemical shifts and coupling constants almost the same, and relaxation times usually not tremendously different. From the activity of 32 Ci/mole, it can be estimated that practical FT measurements on narrow lines will typically employ a few millicuries to a few tens of millicuries of activity in the sample. The soft beta radiation will not penetrate the wall of an NMR tube, making the handling of sealed samples quite safe as long as the tubes remain intact. Reservations about the practicality of ^3H NMR may be at least partially laid to rest by the recent papers from Surrey listed below, as well as by some preliminary work done here.

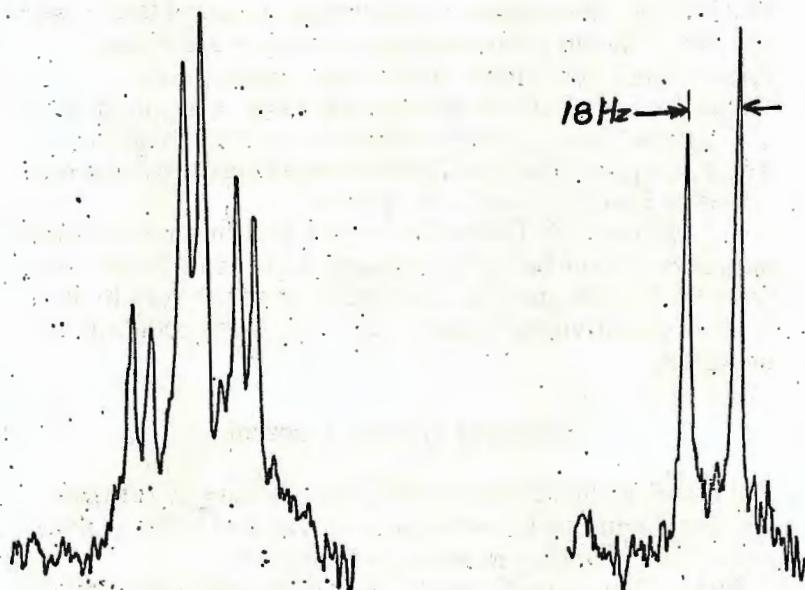
The results of one of our first experiments are shown below. A sample of 50 mCi of "uracil-5- ^3H " in 200 μl of 10% D_2O /EtOH was examined. The uncoupled spectrum on the left and the proton noise decoupled spectrum on the right were found. The only reasonable explanation seems to be the radiation induced hydration:

Professor B. L. Shapiro

August 12, 1975



The usefulness of proton decoupling is obvious in this example, which could otherwise have been rather misleading, and should make possible a number of interesting experiments.



Studies by NMR spectroscopy of stable nuclei in samples containing various other radioactive isotopes might also be safely accommodated, and we would be interested in learning of any needs for such experiments that might exist.

Paul C. Lauterbur
Professor of Chemistry

Yours truly,

Lawrence J. Altman
Assoc. Prof. of Chemistry

1. J. Bloxsidge, J. A. Elvidge, J. R. Jones, and E. A. Evans, *Org. Magn. Res.* **3**, 127 (1971).
2. J. M. A. Al-Rawi, J. A. Elvidge, D. K. Jaiswal, J. R. Jones and R. Thomas, *J. C. S. Chem. Comm.* 220 (1974).

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DEPARTMENT OF CHEMISTRY
(213) 746-2780

August 7, 1975

Professor B.L. Shapiro
Department of Chemistry
Texas A&M University
College Station, Texas 77843

Dear Barry:

Title: ^{29}Si - ^{19}F NOE and T_1 for Hexafluorodisilane

We have determined what we believe to be the first observed NOE for ^{29}Si - ^{19}F . A neat sample of Si_2F_6 (b.p. -18°) in a 9 mm medium wall tube was placed inside a standard 12 mm tube containing C_6D_6 and 19.9 MHz spectra were recorded in the F.T. mode on an XL-100 (probe temperature 35°). Spectra were obtained using gated decoupling (NOE mode) of the ^{19}F decoupling frequency. The value we obtained for η is -0.34 . We have also determined the T_1 for ^{29}Si in the same sample to be 33 sec using the standard inversion-recovery sequence.

The theoretical maximum for the NOE (assuming 100% dipole-dipole relaxation) is -2.37 . These data allow a calculation of the contributions to T_1 as follows:

$$T_1^{\text{DD}} = 230 \text{ sec.}$$

$$T_1^{\text{other}} = 39 \text{ sec.}$$

We assume that T_1^{other} is dominated by contributions from spin-rotation, which should be quite effective for a molecule possessing the high symmetry and volatility of Si_2F_6 . Theoretical reckoning of the value for T_1^{DD} may be rather difficult since the Si-F bond distance in Si_2F_6 is not precisely known and may be significantly different from that in SiF_4 .

Sincerely yours,



Kenneth G. Sharp



Kenneth L. Servis

KLS:KGS:cab

Sony Brook

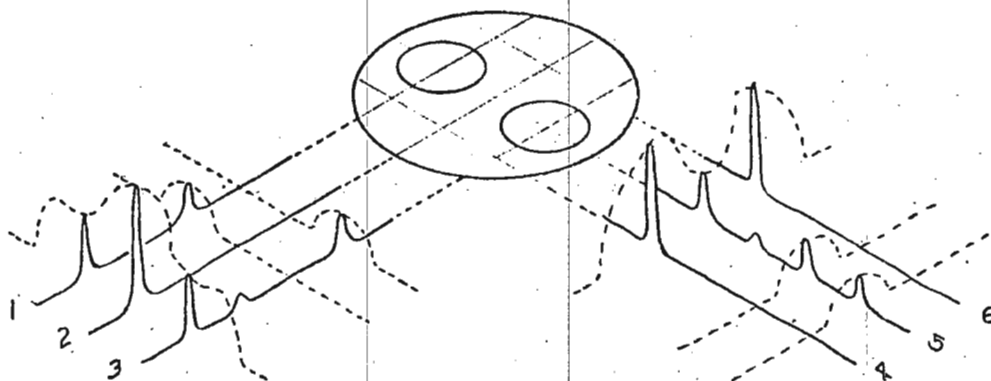
August 12, 1975

Professor B. L. Shapiro
Department of Chemistry
Texas A & M University
College Station, Texas 77843

Dear Barry:

"Zeugmatographic Spectroscopy"

It seems not unlikely that knowledge of the locations within complex objects of various chemical species may at times be of some interest and utility. If their chemical shifts may be distinguished in any NMR experiment in a homogeneous applied magnetic field, their individual distributions within the object may be discovered by combining zeugmatographic spatial discrimination with spectroscopic frequency discrimination. One of the several possible schemes for achieving such a combination may be described with the aid of the diagram below.



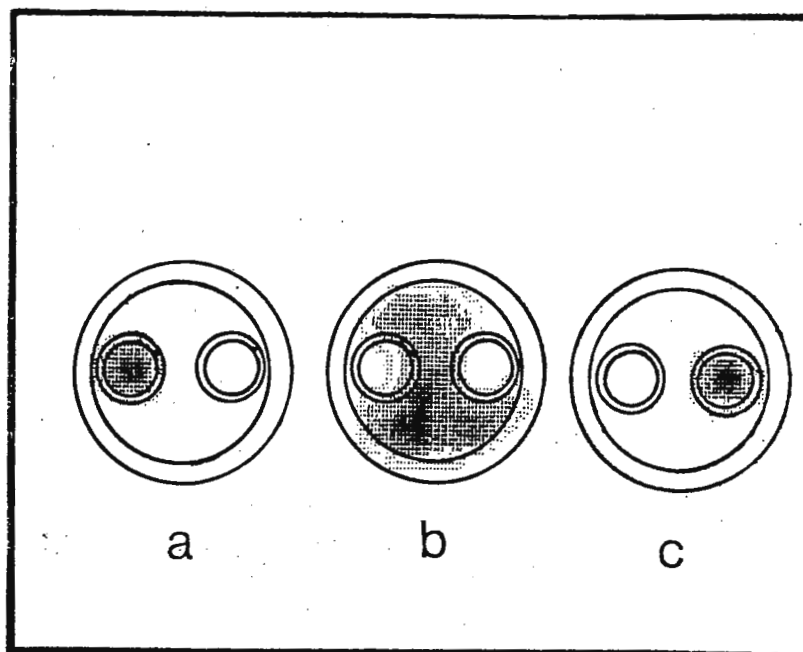
It represents a view along the axis of a set of tubes containing three different compounds in three compartments. In the presence of a field gradient, a long rf pulse may excite only a thin slice of the sample. If the gradient is then switched off, the free induction decay of this selectively excited signal may be observed in a homogeneous field, as in a normal high-resolution pulsed NMR experiment. The Fourier transform of the FID is composed only of the peaks of those compounds in the excited region, as sketched, for example, along line 1 in the diagram. Excitation in another plane with the same gradient

Professor B. L. Shapiro

August 12, 1975

direction may give the spectrum shown as 2, and so on. Separate projections of the intensities of each observable line may be constructed, as shown by the dashed lines. Repetition of the whole procedure for other gradient directions can give other projections, as shown on the right of the diagram.

Finally, by the image reconstruction techniques used in earlier zeugmatographic experiments, the separate distributions of the three compounds may be found. The actual test sample used by Dave Kramer in these experiments had sulfuric acid in the outer tube, water in one of the inner tubes, and $(\text{CH}_3)_3\text{C}-\text{C}_6\text{H}_4-\text{NO}_2$ in the other. The images are shown below, and clearly demonstrate that the objective has been achieved.



Until we get the whole process under computer control, such an experiment, which required special modifications and additions to our Seimco pulsed NMR system by Waylon House, and computer programming and processing by Ching-Nien Chen, is tedious in the extreme. Nevertheless, it may eventually find use with various objects of technological or biological significance from which narrow liquid lines may be obtained, or whose solid resonances may be narrowed by multiple pulse or double resonance methods.

Yours truly,

Paul

Paul C. Lauterbur
Professor of Chemistry

PCL:eg



DEPARTMENT OF HEALTH, EDUCATION, AND WELFARE
PUBLIC HEALTH SERVICE
NATIONAL INSTITUTES OF HEALTH
BETHESDA, MARYLAND 20014

August 12, 1975

Dr. B. L. Shapiro
Department of Chemistry
Texas A & M University
College Station, Texas 77843

Dear Dr. Shapiro:

NMR Study of the Mechanism of Isomerization
of a Series of Aldimines

Recently we have developed a method for enhancing the normally low concentrations of less stable syn isomer in aldimines.¹ The ability to produce the less stable isomer under conditions where its lifetime was much longer than kinetic measurements of its decay has allowed us to study the mechanism of the syn to anti isomerization.

Our preliminary results (see Table) of a series of aldimines as a function of substituents and the medium (solvent effects) are interesting and strongly suggest that isomerization about the C=N bond in these compounds occurs by more than one mechanism. For a rotation mechanism initial bond rupture is heterolytic and therefore the transition state should be more stable in a polar solvent, i. e. rate should be faster in methanol than in benzene (or toluene). In the case of a lateral shift mechanism some redistribution of charged density on the carbon and nitrogen is expected, however the magnitudes should be less. As one can see for compound 1 the rate increased by a factor of 160 in going from toluene to methanol, however, the activation energies remains about the same (7-10 kcal) in both solvents. The result suggests that 1 isomerizes predominantly via rotation mechanism in both solvents. The largest increase in rate about 10^3 occurred for 3, however the activation energies for isomerization in benzene and in methanol differed appreciably indicating a probable change in mechanism (from lateral shift to rotation) in going from a nonpolar to a polar solvent. The rate of isomerization in 4 which has a large activation energy in both solvents actually decreases in going from benzene to methanol, suggesting that 4 proceeds predominantly via a lateral shift mechanism. Compound 2 probably isomerizes by both mechanisms in toluene and methanol.

A complete discussion of these results will be reported shortly.

Sincerely yours,

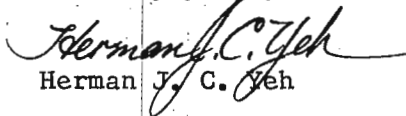
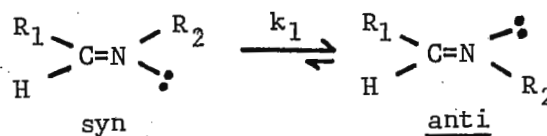

Herman J. C. Yeh

Table. Rates and Activation Energies for the syn → anti Isomerization in Aldimines.



	Compound	Solvent	k_1 , sec ⁻¹ at 30°	E _a kcal/mole
<u>1</u>	R ₁ = 1-naphthyl	Toluene-d ₈	5.5x10 ⁻⁵	7.0
	R ₂ = methyl	Methanol-d ₄	9.0x10 ⁻³	10.2
<u>2</u>	R ₁ = 1-naphthyl	Toluene-d ₈	4.0x10 ⁻³	17.0
	R ₂ = <u>tert</u> -butyl	Methanol-d ₄	2.0x10 ⁻²	12.8
<u>3</u>	R ₁ = 9-anthryl	Benzene-d ₆	6.0x10 ⁻⁸	26.2
	R ₂ = methyl	Methanol-d ₄	6.0x10 ⁻⁵	10.7
<u>4</u>	R ₁ = 9-anthryl	Benzene-d ₆	3.2x10 ⁻⁴	29.5
	R ₂ = <u>tert</u> -butyl	Methanol-d ₄	8.0x10 ⁻⁵	21.3

1. H. J. C. Yeh, H. Ziffer, D. M. Jerina and D. R. Boyd, J Amer. Chem. Soc., 95, 2741 (1973)

CALIFORNIA INSTITUTE OF TECHNOLOGY

PASADENA, CALIFORNIA 91109

August 13, 1975

DIVISION OF CHEMISTRY AND CHEMICAL ENGINEERING
GATES AND CRELLIN LABORATORIES OF CHEMISTRYJOHN D. ROBERTS
INSTITUTE PROFESSOR OF CHEMISTRY

Professor B. L. Shapiro
Department of Chemistry
Texas A and M University
College Station, Texas 77843

Natural-Abundance ^{15}N Nmr
with the WH-180

Dear Barry,

Many of your readers heard the account by Devens Gust at the ENC meeting in April of the good works possible with the WH-180 of natural-abundance ^{15}N nmr spectra of rather large molecules. This does not mean there are no problems; in fact, there are several, the most important being the rather difficultly predictable NOE of ^{15}N and the very long relaxation times associated with tertiary amines, as in ordinary tertiary amines.

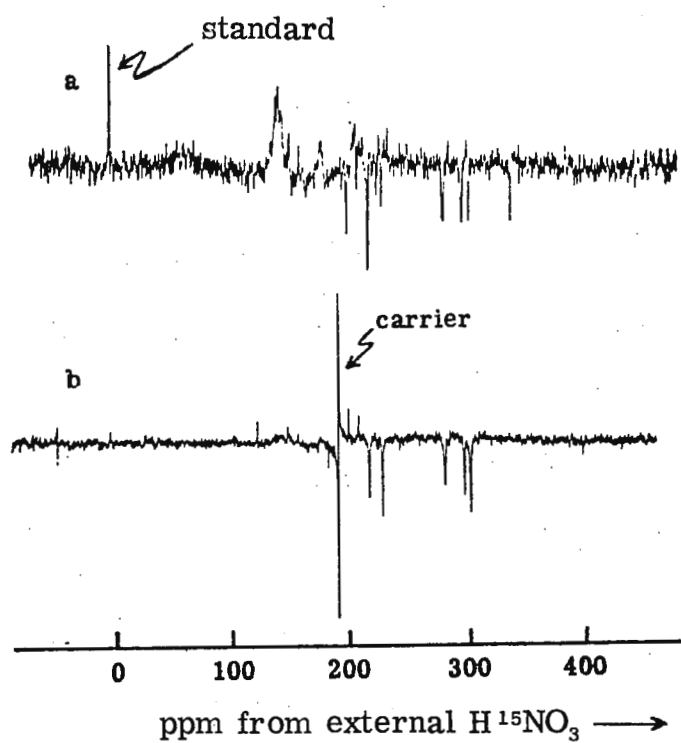
Six years ago, we had real difficulty observing ^{13}C at 0.15 M concentrations and, because ^{15}N at natural abundance is expected to take about a 6×10^{11} longer observation period for a comparable signal-to-noise ratio, we were not optimistic about getting useful ^{15}N at similar concentrations. We still find it hard to believe that even with the larger tubes (25 mm), higher field (45 kgauss), and quadrature detection of the WH-180, combined with the shorter relaxation times arising from longer correlation times with large molecules, have permitted us to take spectra down to 0.015 M or less with natural-abundance ^{15}N . An excellent example is provided by spectra taken here by Richard Moon and Devens Gust of yeast tRNA as shown below. To be sure, an outlandish quantity of tRNA was required—but information is there, and in reasonable amount. The differences between the helical structure and the "melted" form coming with a temperature differences of about 50° are quite striking. The resonances of tRNA at 80° could be assigned to nitrogens of guanosine, cytidine, adenosine and uridine by comparison with the chemical shifts of the 5'-monophosphates measured here by Volker Markowski.

With all good wishes,

Very truly yours,

Jack

JDR:k



Proton decoupled spectra of yeast tRNA (5 g/15 ml solution, about 11 mM) in 0.15 M sodium chloride solution. Spectrum (a) was obtained on a sample at pH 5.6 at about 30° with a 90° pulse, a 0.409 sec repetition rate, and 301.802 transients. Spectrum (b) was obtained at about 80° on a sample at pH 5.4 using a 0.819 sec repetition rate and 82,810 transients.

UNIVERSITY OF CALIFORNIA, DAVIS

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DEPARTMENT OF CHEMISTRY

DAVIS, CALIFORNIA 95616

August 15, 1975

Professor B. L. Shapiro
 Department of Chemistry
 Texas A & M University
 College Station, Texas 77843

Shift-Reagent-Induced Relaxation

Dear Barry,

We have been interested for some time in the use of lanthanide-induced paramagnetic relaxation in organic substrates as a solution structural probe to complement shift reagent studies. Although relative linewidths have been used to estimate relative values of r^{-6} , there are several reasons why such T_2 data cannot be simply interpreted in terms of substrate structure. In addition to the fact that the relaxation rates in a magnetically anisotropic system depend on the structure in a more complicated manner than just r^{-6} , (Texas A & M Newsletter #183, p. 6) sizable contributions from chemical exchange effects can and do occur.

The relative induced linewidths for a substrate, (S), proton have generally been taken as indicative of the inherent relaxation rates induced by different lanthanide complexes, (SR). Such linewidth, δ , as well as shift, ΔH , data, for o-H of 3,5-lutidine coordinated to some SR's in CCl_4 , ($[SR]/[S] = 0.1$), are given in the Table. A better indication of the inherent spin-lattice relaxation rates due to different SR's, however, can be obtained from the protons on the SR-complexes, (i.e. methine proton, β -H, in $Ln(dpm)_3$). These β -H linewidths (as well as shifts), are insensitive to the $[SR]/[S]$ ratio as long as S is in large excess. These β -H shift and linewidths are also listed in the Table.

Some conclusions indicated by the data in the Table are: 1) the β -H shift is primarily dipolar, since $\Delta H(o-H)/\Delta H(\beta-H)$ is essentially invariant with Ln; 2) the inherent relaxation rates indicated by the different lanthanides do not differ very much and tend to parallel the relative values of μ_{eff}^2 ; 3) the o-H linewidths exhibit large exchange broadening contribu-

Professor B. L. Shapiro

August 15, 1975

tions for some Ln since the relative $\delta(\text{o-H})$ do not parallel $\delta(\beta\text{-H})$, but correlate better with the magnitude of the induced shifts. Some current T_1 measurements should provide more definitive answers to this problem.

Sincerely,



Gerd N. La Mar
Professor of Chemistry

GNL:gmh

TABLE

<u>Ln(dpm)₃</u>	<u>3,5-Lutidine^a</u>		<u>dpm</u>		
	<u>$\Delta H(\text{o-H})^b$</u>	<u>$\delta(\text{o-H})^e$</u>	<u>$\Delta H(\beta\text{-H})^b$</u>	<u>$\delta(\beta\text{-H})^c$</u>	<u>$\Delta H(\text{o-H})/\Delta H(\beta\text{-H})$</u>
Tb(dpm) ₃	+ 41	180	- 135	102	- 0.30
Dy(dpm) ₃	+ 53	270	- 159	112	- 0.33
Ho(dpm) ₃	+ 24	160	- 79	92	- 0.30
Er(dpm) ₃	- 11	70	+ 37	54	- 0.30
Tm(dpm) ₃	- 25	150	+ 83	26	- 0.30

a) $[3,5\text{-Lutidine}]/[\text{Ln(dpm)}_3] = 0.10$, in CCl_4 solution at 35°C .

b) Shift, in ppm, from same solution containing La(dpm)_3 .

c) Linewidth, in Hz at 100 MHz.



UNIVERSITY OF SOUTH FLORIDA

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DEPARTMENT OF CHEMISTRY
TAMPA, FLORIDA 33620

813: 974-2571

August 16, 1975

Prof. B. L. Shapiro
Department of Chemistry
Texas A&M University
College Station, TX 77843

Dear Barry:

It is now well established that interactions between lanthanide shift reagents (LSR) and organic Lewis bases can be quite well accounted for by the simultaneous presence of two complexed species: LS and LS₂. The possible implications of the presence of two complexes on structure assessments have, however, been ignored by many workers. In this short letter we show a case in which different bond lengths between the LSR metal atom and the substrate basic substituent atom are obtained when different measures of the lanthanide-induced shift (LIS) are used.

The substrate in this case is 1-adamantanecarbonitrile (1-cyano-adamantane) and the LSR is Eu(fod)₃. Similar behavior has been observed by us for Pr(fod)₃ and for several other nitriles. The shift parameters investigated are the bound shifts of the LS and LS₂ species (Δ_1 and Δ_2), the initial slopes of the LIS dilution curves (λ), and isolated observed shifts at $\rho = 0.5, 1.0, 1.5$, and 2.0 . These LIS values were fit to the simplest form of the pseudocontact equation:

$$S_i = \frac{k(3\cos^2\theta_i - 1)}{r_i^3}$$

Here, S_i simply represents the particular shift parameter under consideration. Since r_i and θ_i are not independent variables for any particular (or rigid) structure, only two parameters enter into the fits: k and the N-Eu bond distance (R_{LN}).

The results are presented in the accompanying table. It is immediately apparent that the values derived for k and the bond length are very much dependent on the particular LIS parameters used. Of course, many more sets of data must be gathered and correlated before a definitive set of conclusions can be reached. For instance, here we are fitting two parameters to only four different protons in a single substrate. Nevertheless, it now appears that the simultaneous presence of LS and LS₂ necessitates the use of bound shifts.

For the sake of brevity, many numbers are not included in the table. Only the results, not the data from which they were derived, are reported. Anyone wishing a closer look at these latter items

is certainly welcome to write to one of us.

Sincerely yours,

Milt
Milton D. Johnston, Jr.
Assistant Professor

Young
Douglas J. Raber
Associate Professor

TABLE

STRUCTURE FIT RESULTS FOR 1-ADAMANTANECARBONITRILE*

Parameter	$R_{LN}^{\circ}(\text{\AA})$	$10^{-2}k$	R
Δ_1	1.98	8.28	0.995
Δ_2	1.71	3.58	0.996
λ	1.67	7.34	0.996
$\rho = 0.5$	1.79	3.62	0.999
$\rho = 1.0$	1.96	6.08	0.997
$\rho = 1.5$	1.94	6.85	0.996
$\rho = 2.0$	1.93	7.19	0.996

*Data were obtained by methods described elsewhere (B. L. Shapiro and M. D. Johnston, Jr., J. Amer. Chem. Soc. 94, 8185 (1972); M. D. Johnston, Jr., B. L. Shapiro, et al., J. Amer. Chem. Soc. 97, 542 (1975)). The k values have units of ppm. The "R" in the last column is the correlation coefficient (subtract it from unity and you get the so-called "agreement factor."); differences in "R" from result to result are not statistically significant. Note that all these parameters afford "good" fits to the structure--at least under the criterion of the correlation coefficient--but clearly only one value of R_{LN} can be correct. Of all the criteria one can use for goodness-of-fit, it can be plainly seen here that the correlation coefficient can be quite insensitive.



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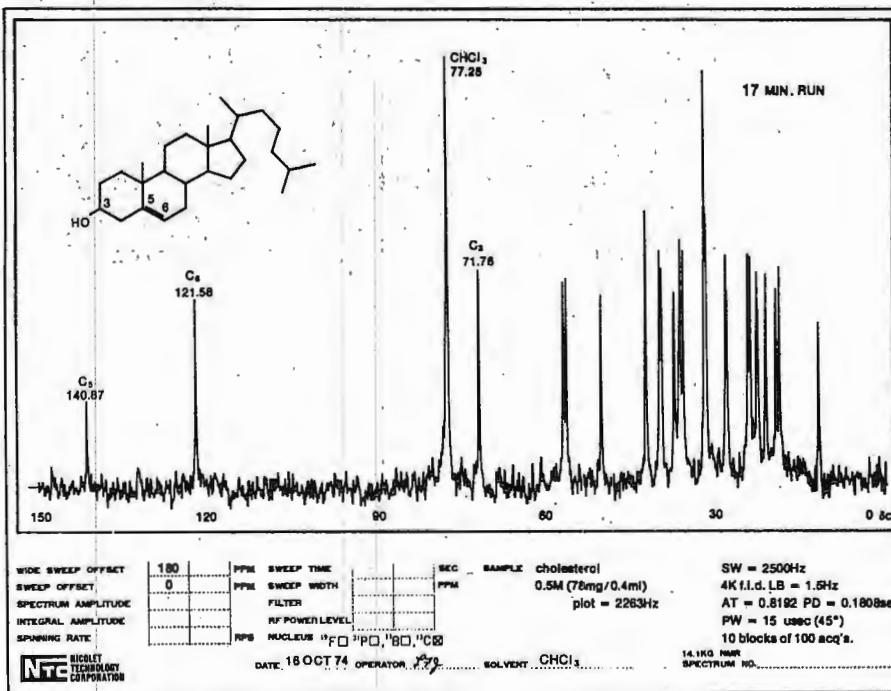
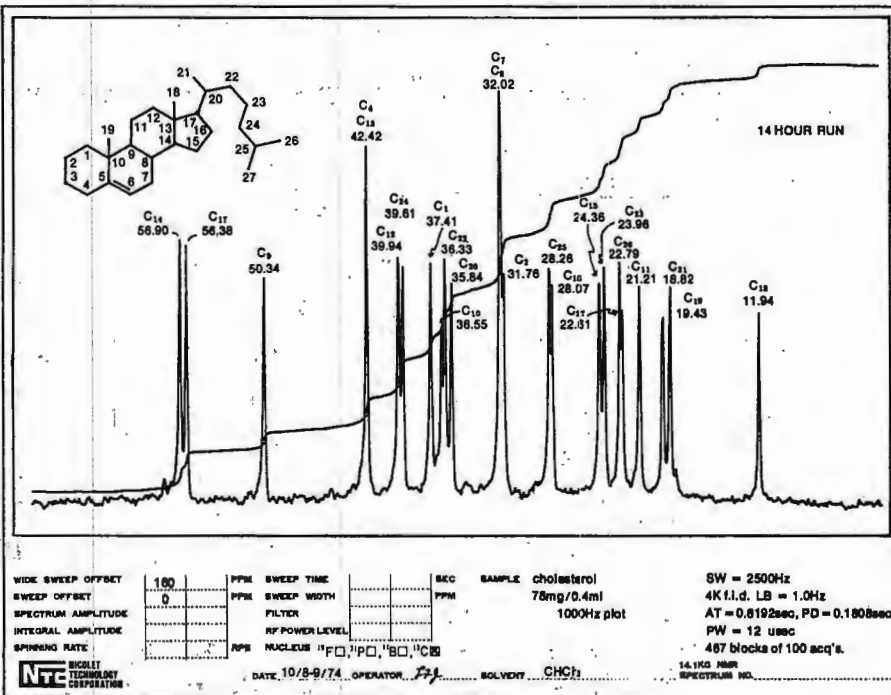
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CENTRAL RESEARCH

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203-445-5611

August 8, 1975

B. L. Shapiro
Dept. of Chemistry
Texas A & M University
College Station, TX 77843

Title: Random Noise Stirring

Dear Barry:

Having also rediscovered Ernst's noise modulation experiment at intermediate RF power levels, we were very interested in Hall's recent letter on the subject (202-24). Since the conditions are rather analogous to those described by Baldeschweiler (for coherent radiation) as "stirring", and mnemonics are in vogue these days, we refer to the experiment as "random noise stirring" (RNS). As an alternative, "agitation" might be an equally appropriate term. We could then distinguish between spatially coherent, decoupler noise (DNA) and truly random (RNA) processes, and confuse readers to the same extent we do with off-resonance decoupling (ORD) experiments.

The experiment has been applied by Schaffer (Macromolecules 5, 590 (1970)) to distinguish methylene from methine and methyl carbons, and it is in this direction that our own interest lies. This is prompted by the fact that the experiment is quite sensitive to intersystem crossing between the singlet and triplet states of the methylene protons, and less to other complications than alternative ways (ORD) of looking for methylene non-equivalences. Hence, it is useful in distinguishing methylene carbons bearing equivalent and nonequivalent protons in the assignment of CMR spectra. As an example, the proton nonequivalence in acetaldehyde diethyl acetal can be detected by the broadening of the methylene carbon line.

We haven't been able to distinguish magnetic equivalence (commutation of F^2) from symmetric equivalence yet, although this would be fun to do. One complication is that proton relaxation can give intersystem crossing and a T_2 effect on the CMR of the singlet, which is not apparent when the system is completely decoupled.

Another fun experiment is to look at the T_1 and NOE of the singlet and triplet combinations separately in magnetically equivalent cases. To play, one should try to guess the outcome before doing the experiment.

Sincerely yours,

E. B. Whipple

/kce

National Research Council
CanadaConseil national de recherches
CanadaDivision of Biological
SciencesDivision des sciences
biologiques

File Référence

12 August 1975

A CONVENIENT MEANS TO REMOVE THE INFLUENCE OF RESIDUAL
PARAMAGNETIC METAL IONS IN DEUTERIUM OXIDE ON
 ^{13}C SPIN-LATTICE RELAXATION TIMES

Dear Barry,

Roberts and coworkers¹ recently reported the effect of traces of metal ions, present as impurity in commercially available D_2O , on the carbon-13 spin-lattice relaxation time (T_1) of the carbonyl carbon of glycine. According to these authors the presence of paramagnetic ions causes a severe shortening of the experimentally determined T_1 values; the effect of these ions can be eliminated by extensive treatment of glassware and sample with EDTA.

We have tried to achieve similar results in a more simple fashion by adding hydrogen sulfide and/or EDTA directly to the sample. Both of these reagents are known to remove metal ions from solutions very efficiently and direct addition to the sample could prove to be a fast and convenient method for removing metallic impurities from commercially available D_2O . Hydrogen sulfide is a particularly attractive reagent as products are easily removed from the sample. Glycine- $1\text{-}^{13}\text{C}$ (Merck, Sharpe and Dohme, Canada) was dissolved in commercial D_2O , the major source of which is the U.S. Atomic Energy Commission. Hydrogen sulfide was bubbled through the solution for about 10 minutes and the T_1 values were determined before and after bubbling. Different samples of glycine were prepared with added paramagnetic ions and T_1 values were measured before and after the addition of hydrogen sulfide or EDTA solution. A Varian CFT-20 spectrometer was used to determine T_1 values via a $(180-\tau-90-T)_n$ sequence.

The data in Table I show that the common impurities in the commercially available D_2O are effectively removed by hydrogen sulfide. After separation from the precipitate, residual H_2S can be removed from the sample by warming under reduced pressure. The T_1 values obtained by this procedure are very close to those reported by Roberts and coworkers¹. This indicates that the method is rapid, convenient, and effective for removing the paramagnetic ions which might affect the spin-lattice relaxation times in D_2O . No extensive treatment of glassware or sample is necessary.

In order to check whether the observed increase in the spin-lattice relaxation time caused by bubbling hydrogen sulfide is due to the partial removal of dissolved oxygen, nitrogen was bubbled through the sample for about 10 minutes before determining the value of T_1 . In a subsequent experiment the nitrogen-saturated sample was subjected to several freeze-pump-thaw cycles and the T_1 values measured. The results in Table I indicate that the increase observed on saturation with H_2S is entirely due to removal of the metal ions from the solution.

To get an idea of the efficacy of the method for particular metal ions, the experiments were repeated with known concentrations of Mn^{2+} , Fe^{3+} , Cr^{3+} , and Cu^{2+} . From Table I it is apparent that cupric ions are completely removed by hydrogen sulfide, ferric ions are removed to a large extent, and manganous and chromic ions are removed only partially by this method. This is in agreement with the solubility products and the solubilities of the corresponding sulfides in slightly acidic solution^{2,3}. Cupric sulfide, having a very low solubility product (4×10^{-38}) at room temperature, is completely precipitated. The sulfides of the other ions, being more soluble in acid, are not fully precipitated. In this case addition of 0.001M EDTA is sufficient to remove the influence of the metal ions.

Best regards,

F. Hasan *Roxanne* *Ian*
Fariza Hasan, Roxanne Deslauriers and Ian C. P. Smith

1. H. Pearson, G. Gust, I.M. Armitage, H. Huber, J.D. Roberts, R.E. Stark, R.R. Vold and R.L. Vold, Proc. Natl. Acad. Sci. U.S., 72, 1599 (1975).
2. E.B. Kelsey and H.G. Dietrich in "Fundamentals of Semimicro Qualitative Analysis", The MacMillan Co., New York, 1956.
3. "Handbook of Chemistry and Physics", The Chemical Rubber Co., 50th ed., 1969-70.

Effects of paramagnetic ions on ^{13}C -spin-lattice relaxation times for the carboxyl carbon of 0.34 M glycine in D_2O .

Added Ions	pD	T_1 (sec)
-	6.9	19.9
with H_2S	5.4(6.4) ^a	83.5
with N_2	6.9	17.9
after degassing	6.9	17.7
Mn^{2+} , 9×10^{-6} M	6.9	7.9
with H_2S	5.5(6.5) ^a	42.0
with EDTA	6.8	75.7
Fe^{3+} , 9.6×10^{-5} M	6.8	6.5
with H_2S	5.5(6.5) ^a	62.7
Cr^{3+} , 1×10^{-4} M	6.9	41.8
with H_2S	5.5(6.5) ^a	47.6
Fe^{3+} , 9.6×10^{-5} M + Cr^{3+} , 1×10^{-4} M		
with EDTA	6.7	75.0
Cu^{2+} , 2.5×10^{-5} M	6.9	9.7
with H_2S	5.4(6.5) ^a	82.5

^apH of the sample after removing hydrogen sulfide.

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DEPARTMENT OF CHEMISTRY
SANTA BARBARA, CALIFORNIA 93106

August 18, 1975

Dr. B. L. Shapiro
Department of Chemistry
Texas A & M University
College Station, Texas 77843

"Surplus Tubes for Sale"

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Sincerely yours,

J. Thomas Gerig
Associate Professor

JTG:jh



THE UNIVERSITY OF WINNIPEG
WINNIPEG, CANADA R3B 2E9

Department of Chemistry
18 August, 1975

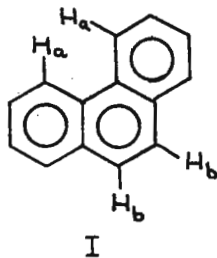
Professor B.L. Shapiro
Department of Chemistry
Texas A and M University
College Station, Texas 77843

Dear Barry:

$^1\text{H}-T_1$'s in condensed ring hydrocarbons

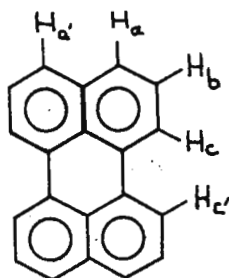
We have been interested in examining the potential of ^1H nmr spin-lattice relaxation rates in studying condensed ring hydrocarbons. Our results indicate that proton T_1 's may be helpful in confirming the assignment of ^1H nmr spectra and in some cases may also allow crude estimates of internuclear proton-proton separations.

For phenanthracene (I) we find $T_1(\text{H}_a) \approx 5.7$ sec and $T_1(\text{H}_b) \approx 11.6$ sec. Since $T_{1,d-d}^{-1}$ depends upon r^{-6} the two

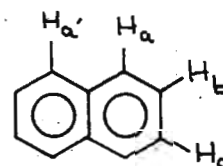


equivalent H_a protons have the shorter T_1 because of their relatively small internuclear separation. In this particular example proton-proton NOE experiments would not be helpful in allowing distinction of the H_a resonances from other proton resonances.

Similar results have been obtained for perylene (II) where



II



III

$T_1(H_a) = 9.2 \pm 1.0$ sec and $T_1(H_c) = 4.2 \pm 0.5$ sec. It is possible to estimate $r(H_c-H_c')$ in II if one makes a number of assumptions, including: (1) the H_a-H_a' , H_a-H_b , and H_b-H_c separations in II are the same as those in naphthalene (III): approximately 2.45 Å, 2.50 Å, and 2.45 Å, respectively, (2) intramolecular dipole-dipole relaxation completely dominates the spin-lattice relaxation of H_a and H_c , and (3) T_1^{-1} is given by an expression similar to that of Gutowsky and Holm [1], $T_1^{-1} = 3/2 \gamma_H^4 \hbar^2 \gamma_c \sum r^{-6}$. We calculate approximately 2.0 Å for $r(H_c-H_c')$ which is in fair agreement with the value estimated from X-ray crystallographic data [2].

Sincerely,

Roderick Wasylishen

Roderick Wasylishen

- [1] H.S. Gutowsky and D.E. Woessner. Phys. Rev. 104, 843 (1956).
- [2] A. Camerman and J. Trotter. Proc. Roy. Soc. (London), 279A, 129 (1964).

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casella postale 3698 . 20100 Milano

destinatario:

Professor
B.L. Shapiro
Department of Chemistry
Texas A and M University
College Station, TX 77843, USA

nostro riferimento Department of Physical Chemistry na: August, 28, 1975

Subject: ¹³C Spin lattice relaxation times of the antibiotic rifampin

Dear Prof. Shapiro,

in the framework of our studies on ansamycin antibiotics, we have run the ¹³C NMR spectrum of rifampin and performed an almost complete assignment of its 43 C atoms. To try to understand about molecular motion, we have recently measured the ¹³C T₁'s by the PRFT technique, using our new Bruker WP-60 instrument. The preliminary results reported in the Table seem to indicate that the whole molecule has a fairly slow rotational reorientation.

The CH₃ groups of the ansa chain, which show longer T₁'s than the backbone carbons, may have shorter effective correlation times, as observed for the angular CH₃ groups of steroids ¹⁾.

We intend to obtain new T₁ data in different solvents and on the two separated moieties of the molecule, i.e., the naphthoquinone chromophore and the ansa chain, to substantiate this hypothesis.

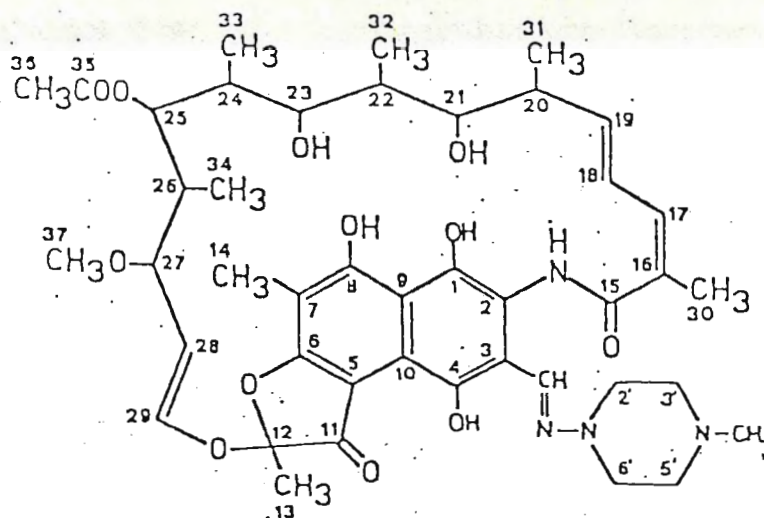
Yours sincerely,

(Edoardo Martinelli)

(Ambrogio Ripamonti)

Edoardo Martinelli *Ripamonti Ambrogio*

1) A.Allerhand, D.Doddrell and R.Komoroski, J.Chem.Phys., 55, 189 (1971)



RIFAMPIN

TABLE - ^{13}C Chemical shifts (δ , ppm) and T_1 's (sec) of rifampin
(conc. ~ 0.3 M in CDCl_3)

C	δ	T_1	C	δ	T_1
1	138.6	3.5	23	76.7	2.1
2	110.8 *	3.5	24	37.6	0.8
3	105.9 *	5.0	25	74.4	0.6
4	147.8	3.3	26	39.5	1.6
5	117.8 *	2.6	27	76.7	2.1
6	174.3	5.5	28	118.7	0.5
7	120.3 *	6.5	29	142.6	1.3
8	169.3	3.2	30	20.7	1.5
9	104.4 *	5.5	31	17.8	1.1
10	112.8 *	6.4	32	10.9	1.2
11	195.3	4.1	33	8.5	0.7
12	108.7	5.2	34	8.8	0.7
13	21.5	1.2	35	171.9	5.3
14	7.6	1.1	36	20.7	1.5
15	169.6	3.2	37	57.0	1.6
16	129.4	3.3	CH=N	134.4	0.4
17	135.0	0.4	2'	50.2	0.6
18	123.2	1.1	3'	53.9	0.6
19	142.6	1.3	5'	53.9	0.6
20	38.6	1.4	6'	50.2	0.6
21	70.7	1.6	N-CH ₃	45.8	1.1
22	33.4	0.7			

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DEPARTMENT OF CHEMISTRY

August 13th, 1975.

Professor Barry L. Shapiro,
 Department of Chemistry,
 Texas A and M University,
 College Station, Texas 77843

Dear Professor Shapiro:

COMPUTATIONALLY OFFSETTING THE FREQUENCY IN NMR SPECTRA
 AND OTHER NMR SIGNALS

It is interesting (and in very limited cases useful) to note that the phase shift routine in fourier transform magnetic resonance spectrometers can be used to computationally offset the frequency. Suppose, for example, that one has obtained an fid with the bad fortune of having one line on exact resonance. Therefore, one would like to shift the entire spectrum by a small number of Hz. The fid is multiplied by $\cos(\omega_s t)$ using the technique described below and then fourier transformed. The line that was on resonance now appears shifted to higher frequencies by ω_s but all lines that were not at zero frequency now appear as doublets at $\omega + \omega_s$ and $\omega - \omega_s$. This is because the process of multiplying by a cosine modulation is identical to mixing the signal with a carrier at ω_s , in an ideal mixer. Owing to folding about zero frequency, the line at zero frequency does not have two sidebands. In order to make the technique useful, one has to remove all lines at $\omega - \omega_s$. This can be done simply: The original fid is multiplied by $\sin(\omega_s t)$ and fourier transformed. The resulting sine transform is identical to the cosine transform in the original shifted spectrum except that now the lower sidebands appear inverted. The sine transform of the second manipulation is then added to the cosine transform from the first manipulation to yield the original spectrum shifted by ω_s . The line originally at zero frequency becomes twice as large in this procedure. Also, ω_s should be chosen so that there were no lines between the zero frequency line and ω_s .

The above procedure is an exact analog of a commonly used electronic scheme for generating variable frequency single-sideband, suppressed-carrier rf in transmitters. The block diagram of a circuit which performs the above procedure is shown.

The procedure for multiplying by a sine or a cosine employs the first-order phase shift routine. As the phase correct routine was developed to operate on spectra after an fid had been fourier transformed, with the cosine (real) transform in the first block and the sine (imaginary) transform in the second block, this routine mixes appropriately data from both blocks in performing the phase correction.

$$X'_{\text{REAL}} = X_{\text{REAL}} \cos\phi + X_{\text{IMAGINARY}} \sin\phi$$

$$X'_{\text{IMAGINARY}} = X_{\text{REAL}} \sin\phi + X_{\text{IMAGINARY}} \cos\phi$$

August 13, 1975.

For our purposes, the fid occupies the region "X_{real}", and X_{IMAGINARY} is zeroed. After "phase correction", the original fid multiplied by $\cos\phi$ is then in the first block and that multiplied by $\sin\phi$ is in the second block. Most first-order phase correction routines cause the last point (the Nyquist frequency point) to be corrected a number of degrees designated by a constant (PB), and the first point left unchanged. All points inbetween are corrected proportionately. The value for PB was determined by taking the acquisition time for the block of data and determining how many cycles of oscillation in that time duration gives rise to the desired frequency offset. This number of cycles $\times 360^\circ/\text{cycle}$ gives the appropriate value of PB.

This technique is of limited utility for shifting ordinary spectra; it might find use in those cases where one must gather the data on resonance and then wishes to fourier transform afterwards (e.g. the Jeener echo experiment).

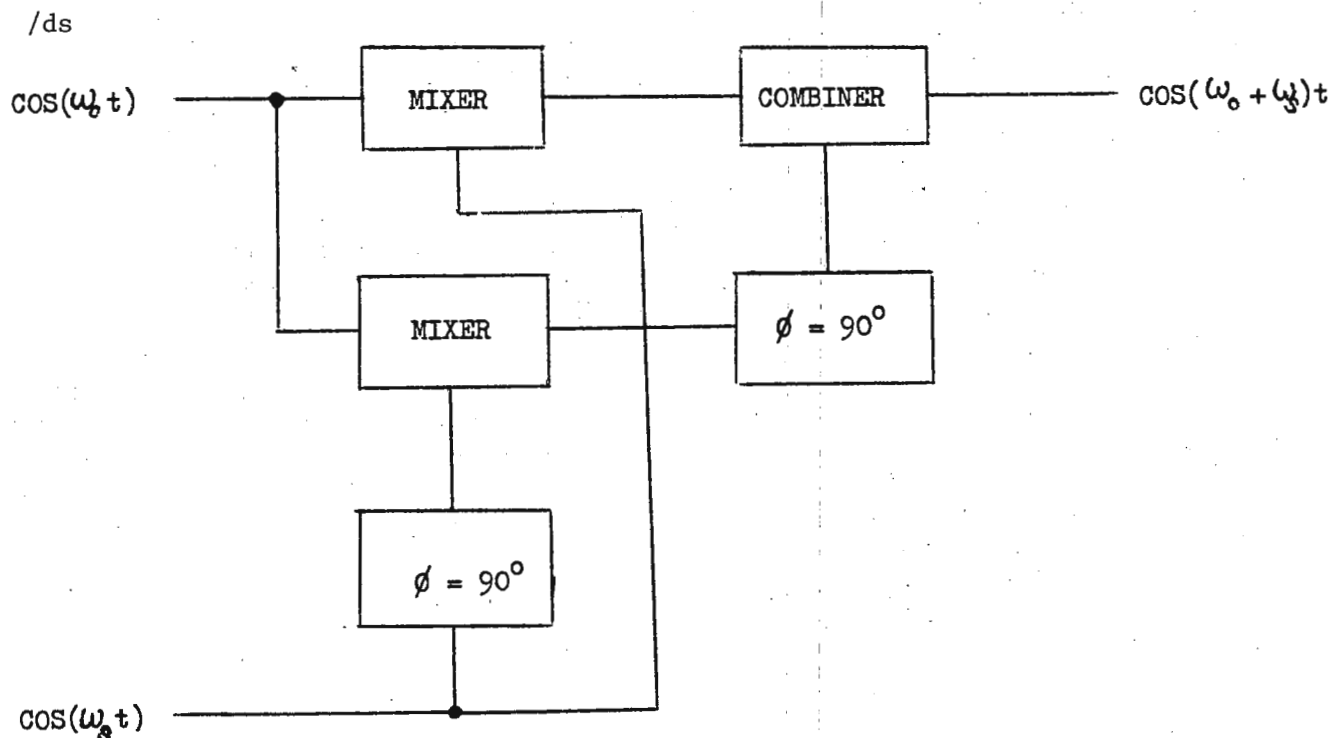
Sincerely,

Stephen B.W. Roeder

Timothy P. Higgs

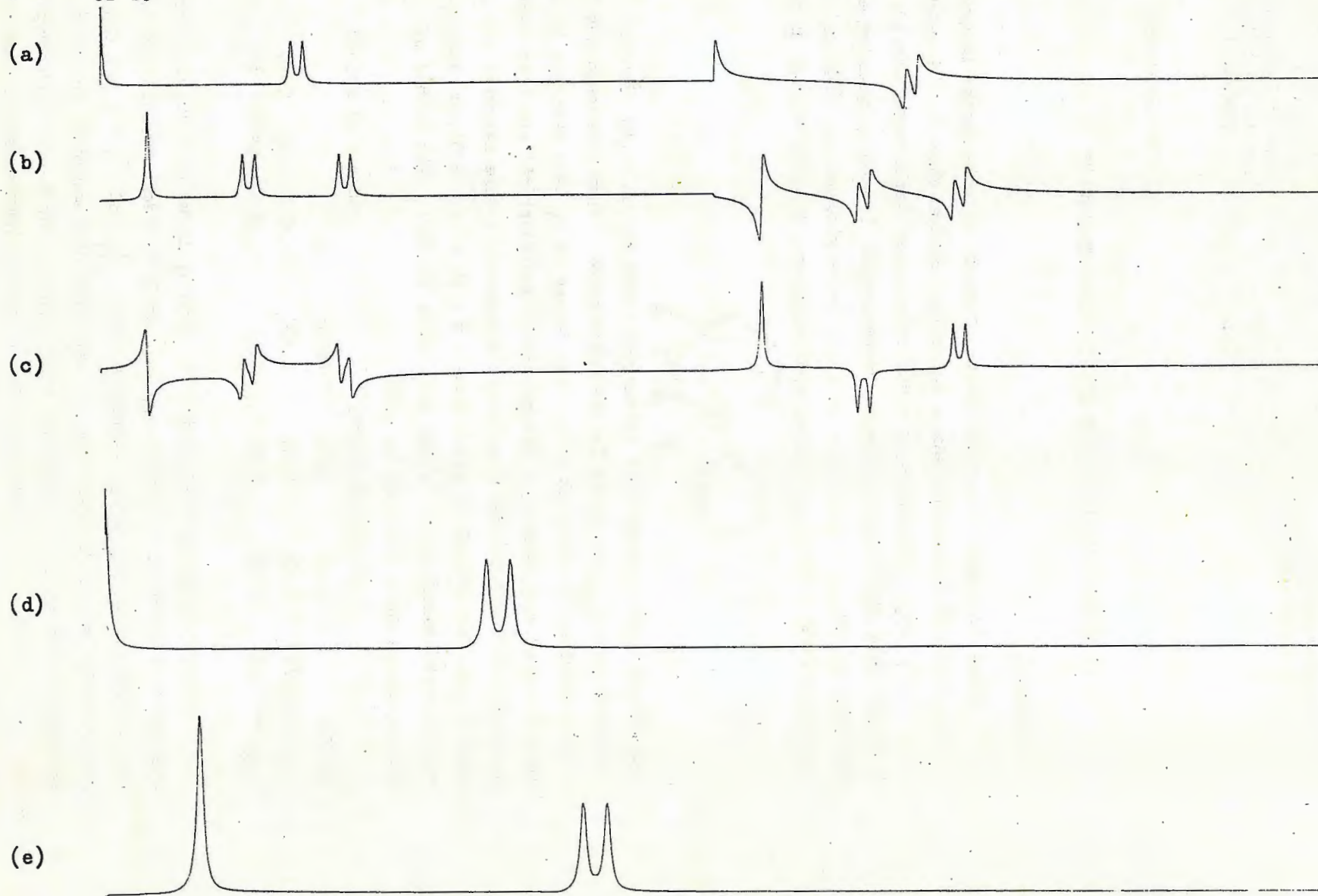
STEPHEN B.W. ROEDER and

TIMOTHY P. HIGGS.



Block diagram of circuit used to generate carrier-suppressed, single-sideband modulation

Top to bottom: (a) Real and imaginary parts of spectrum with one line on exact resonance; (b) Result of modulation of fid by a cosine modulation at frequency ω_s and fourier transformation; (c) Result of modulation of fid by a sine modulation and fourier transformation; (d) Real part of original spectrum; (e) Real part of final spectrum resulting from addition of the right hand side of c to the left hand side of b.



From Dr. R. K. Harris

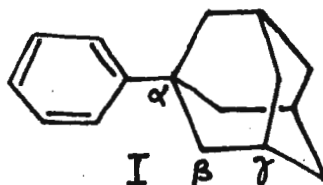
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 Telephone Norwich (0603) 56161
 Telegrams UEANOR NORWICH
 ENGLAND

1st September, 1975

¹³C SPIN-LATTICE RELAXATION FOR 1-PHENYLADAMANTANE

Dear Barry,

Along with many other NMR spectroscopists we have become interested in the detailed information about molecular motion which can be obtained from ¹³C T₁ data. We wished to study relatively simple cases involving a single axis of internal molecular motion which is also a symmetry axis for each moiety. As an example we chose 1-phenyladamantane (I), which has six different types of carbon atom which are directly bonded to protons.



The experimental spin-lattice relaxation times due to (C,H) dipolar interactions, T_{1dd}, are given in the Table below. They were obtained for a 1.0M solution in CDCl₃ at 32°C. The values of T₁ were measured by the inversion-recovery technique (Freeman-Hill modification) and were reproducible to within 3%. The nuclear Overhauser enhancements were measured using gated decoupling; they ranged in value from η = 1.89 to η = 2.01, and were mostly within experimental error of the full value (1.988). The values of T_{1dd} are estimated to be accurate to ± 4%.

	phenyl moiety			adamantyl moiety		
Carbon	para	meta	ortho	β	γ	δ
T _{1dd} (expt)/s	2.89	6.28	6.50	4.32	7.35	2.30
T _{1dd} (calc)/s	2.84	6.32	6.44	4.23	7.43	2.38

We have attempted to describe the data in terms of a single correlation time for end-over-end rotation, τ₁, and two independent correlation times for rotation about the unique molecular axis, τ_{//}^p and τ_{//}^a for the phenyl and adamantyl moieties respectively, using standard geometry information. We have not attempted to correlate the motion of one moiety with respect to the other. We have, in effect, used Woessner's equation for anisotropic rotation¹ of a single rigid molecule, and applied it to each moiety of I

independently. The best fit for the six items of experimental information is given when the correlation time parameters have the following values:

$$\tau_{\perp} = 15.6 \text{ ps} \quad \tau_{//}^p = 2.7 \text{ ps} \quad \tau_{//}^a = 2.3 \text{ ps}$$

As expected τ_{\perp} is considerably longer than $\tau_{//}^p$ or $\tau_{//}^a$. Somewhat more surprising is the fact that the bulkier adamantyl portion appears to rotate somewhat faster about the internal rotation axis than does the phenyl moiety.

The values of T_{1dd} calculated from the above three parameters are given in the Table. It can be seen that they agree with the observed values to within experimental error, thus showing that the three-parameter model is adequate in practice, at least in this case. In the calculations all non-bonded protons on the same moiety as the carbon in question have been taken into account. If only directly-bonded protons had been considered, the ratio T_{1dd}^Y/T_{1dd}^B (adamantyl portion) would have been 1.98; inclusion of non-bonded protons reduced this ratio to 1.75, compared to the experimental ratio of 1.70. Of course, interactions between protons on one moiety and carbons in the other would require a further motional parameter, but it looks as though the data obtained would not justify inclusion of this extra complication, and we have reason to believe that it can be neglected in the present case.

We hope this report is of interest to some TAMUNMR readers, but in any event we trust it suffices to keep RKH on the TAMUNMR mailing list.

Best wishes,

Robin

Roger

R. K. Harris

R. H. Newman

¹ D. E. Woessner, J. Chem. Phys. 37, 647 (1962).

Professor B. L. Shapiro,
Department of Chemistry,
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U.S.A.

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Texas A. & M University
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Ihr Zeichen

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VW/DNL/uh

Tel. Durchwahl:
7008 362/363

Datum
September 3rd, 1975

Measuring accurate line positions on CFT-20

Dear Barry,

We have previously reported¹ an improved way to measure line positions on our XL-100, using an interpolation routine. Recently we needed to perform high-resolution work on our CFT-20 and found that the current version of the software FT-16/T₁ (Varian part No. 994114-07 Rev. E) gave larger r.m.s. errors, at small spectral widths, than we had expected from the Hz/channel ratio.

Perusal of the software showed that in the routine CFIT the frequency correction for point interpolation is being properly calculated, i.e. in channel units, but not in the form that it is used later in the program, i.e. in Hz. Also, the correction is added, at the point of application, to the frequency of the highest-intensity point, instead of being subtracted.

We have cured this problem with a patch to CFIT (see below) that returns a correction in Hz, which is added at the appropriate point in the program. At present this patch overwrites part of the Sykes cassette routine.

We have checked out the patch by recording the proton-coupled ¹³C spectrum of pyrazine, which shows a small coupling of 2.11 ± 0.03 Hz (XL 100) that appears six times in the spectrum. The reproducibility under various conditions is shown in the table.

Yours sincerely,

Dave *V. Wray*
(Dr.D.N. Lincoln) (Dr. V. Wray)

PATCH TO CENTROID FOR CFT-20 PROGRAM (REV.E.)

(11630)	12177	PATCH		
(12177)	O	PATCH	ENTR	
(12200)	64015	STB	CORR	
(12201)	5001	TZA		
(12202)	21041	LDB	SWP	WIDTH IN Hz
(12203)	30271	LDX	DTOX	REAL DATA PTS
(12204)	2000	CALL	GDIV	Hz/CHAN
(12205)	1735			
(12206)	5014	TAX		
(12207)	5001	TZA		
(12210)	164005	MUL	CORR	
(12211)	5042	TXB		
(12212)	164003	MUL	CORR	CORR * Hz/CHAN
(12213)	1000	RETU*	PATCH	
(12214)	112177			
(12215)	O	CORR	DATA	O

PATCH FOR TWO DECIMAL-PLACE FREQUENCY PRINTOUT

(11713) 140300

(11770) 150402

	<u>Table</u>	
<u>Coupling</u>	<u>rms error</u>	<u>Hz/channel</u>
* 1.83	0.51	0.073
2.13	0.03	0.073
2.17	0.17	0.292

* without correction

DEPARTMENT OF CHEMISTRY
THE UNIVERSITY OF GEORGIA ATHENS, GEORGIA 30602

August 21, 1975

Professor B. L. Shapiro
Department of Chemistry
Texas A and M University
College Station, Texas 77843

N-Nitroso-N-methylaniline, An Error

Dear Barry:

We have spent some time trying to reproduce the ^{13}C NMR spectrum of N-nitroso-N-methylaniline as reported previously! It had been reported that the ortho and meta carbons were nonequivalent as a result of slow rotation about the carbon-nitrogen bond. In each of our spectra, the ortho and meta carbons remained equivalent. Finally, I contacted Professor Randall concerning this problem. In his reply, he stated that the reported nonequivalence was spurious and due to an error in data handling and suggested that I use TAUMU NMR to alert others to this error. It is my understanding that Professor Stothers has also tried to reproduce the non-equivalence.

The chemical shifts one obtained for neat N-nitroso-N-methylaniline (5000 Hz, d_6 -acetone lock) with respect to TMS are: CH_3 -30.82, C_{ortho} -118.57, C_{para} -126.58, C_{meta} -128.91, and C_{sub} -141.77 ppm.

Sincerely yours,

Dick

Richard H. Cox
Associate Professor

1. P. S. Pregosin and E. W. Randall, Chem. Commun., 399 (1971).

RHC:mjd

UNIVERSITY of PENNSYLVANIA

PHILADELPHIA 19174

The School of Medicine

JOHNSON RESEARCH FOUNDATION G4
DEPARTMENT OF
BIOPHYSICS AND PHYSICAL BIOCHEMISTRY

3 September 1975

Dr. B. L. Shapiro
Department of Chemistry
Texas A & M University
College Station, Texas 77843

Dear Dr. Shapiro:

Rotameric Preference of β -Alanine

In relation to recent studies on the time-averaged conformation of biological molecules, we have had occasion to study β -alanine. The 60 MHz spectrum in D_2O gave rise to the 14 line AA'BB' pattern shown in Figure 1. The upfield multiplet was assigned to the methylene protons adjacent to the carboxyl group by virtue of the $N-C-C-H$ coupling. The 220 MHz spectrum consisted of two triplets with an average peak separation of 6.65 Hz.

The chemical shifts and coupling constants for the 60 MHz spectrum were obtained by computer assisted analysis using LAOCN III as previously described.¹ This procedure provided vicinal coupling constants of 6.29 and 7.26 Hz, with geminal coupling constants of -12.19 and -14.39 Hz and an rms error of 0.08. These values were used to calculate the 220 MHz spectrum, and a good fit was obtained.

The vicinal coupling constants were used to establish the rotameric preference using the relationship between the electronegativities of substituents and the average coupling constant:²

$$17.97 - 0.8 \sum_{i=1}^6 E_i = 1/3 (3/2 N + 1/2 L)$$

The value of $\sum E_i$ calculated with L negative, 14.37, agrees more closely with that obtained using Huggins electronegativities, 14.45, than does that calculated with L positive, 13.57. Based upon the usual arguments, the trans rotamer is thus favored. Using model coupling constants J_t and J_g we calculate the trans:gauche:gauche rotamer ratio to be $\sim 2:1:1$.

The pattern of triplets at 220 MHz arises because $\frac{L^2}{2|M|} < \Delta\nu_{1/2}$, where L and M are the difference between the vicinal and geminal coupling constants and $\Delta\nu_{1/2}$ is the observed width of a single line, and does not indicate that $J = J'$ with rotamer populations identical. This represents but another example of an old caveat.

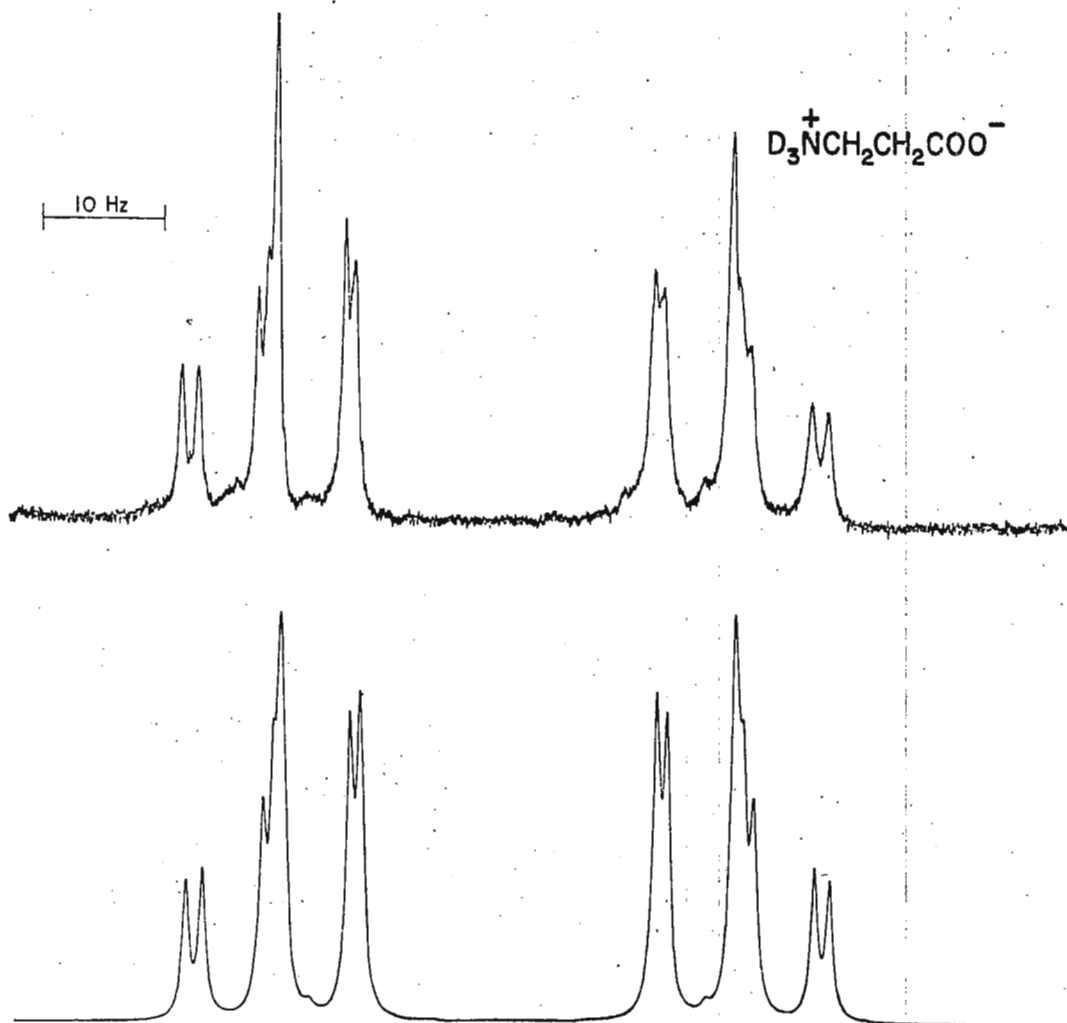
A more detailed report of this work will appear in a forthcoming publication.

1. G.E. Wilson, Jr. and T. J. Bazzone, J. Amer. Chem. Soc., 96, 1465 (1974).
2. R. J. Abraham and K. G. R. Pachler, Mol. Phys., 7, 165 (1963).

Sincerely,

G. Edwin Wilson, Jr.

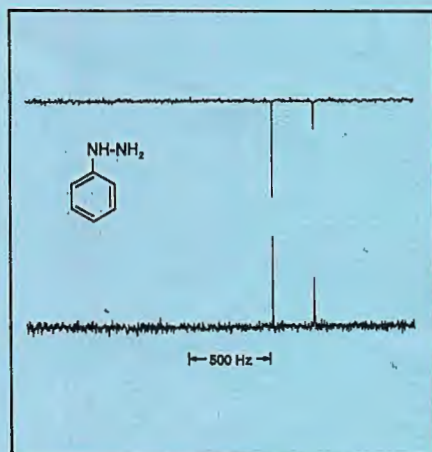
G. Edwin Wilson, Jr.
Visiting Associate Professor



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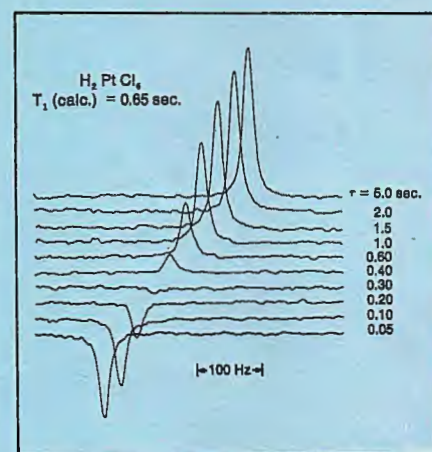


The nucleus observed for this spectrum is ^{15}N , at 10.1 MHz. The upper trace shows 500 transients ($\alpha = 90^\circ$) of a proton noise-decoupled spectrum of phenylhydrazine in C_6D_6 . The negative magnetogyric ratio of ^{15}N produces negative NOE, hence the inverted lines in the trace. The lower trace shows 2000 transients of phenylhydrazine ($\alpha = 90^\circ$); the decoupler was on during acquisition and off during the pulse delay. This technique makes it possible to measure NOE while retaining the advantages of a ^1H noise-decoupled spectrum.

degree of experimental freedom is offered for an NMR Spectrometer that combines state-of-the-art performance and ease of operation. At present, we cannot begin to assess the impact the new-found experimental scope might have on the direction of future investigations. But we expect that a lot of new ground will be broken.

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In this spectrum, the new accessory allows the observation of ^{195}Pt at 21.5 MHz; the sample was aqueous hexachloroplatinic acid. An inversion recovery ($180^\circ - \tau - 90^\circ$) pulse sequence was used in the automatic measurement of the spin-lattice relaxation time (T_1) for the ^{195}Pt nucleus.

We wish to acknowledge the cooperation of Professor Paul Ellis, of the University of South Carolina, whose early experimental work contributed to development of this capability of the XL-100.

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